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# Matrix Models, Topological Strings, & $\mathcal{N}=1$ Gauge Theory

Robert Dijkgraaf  
(Amsterdam)

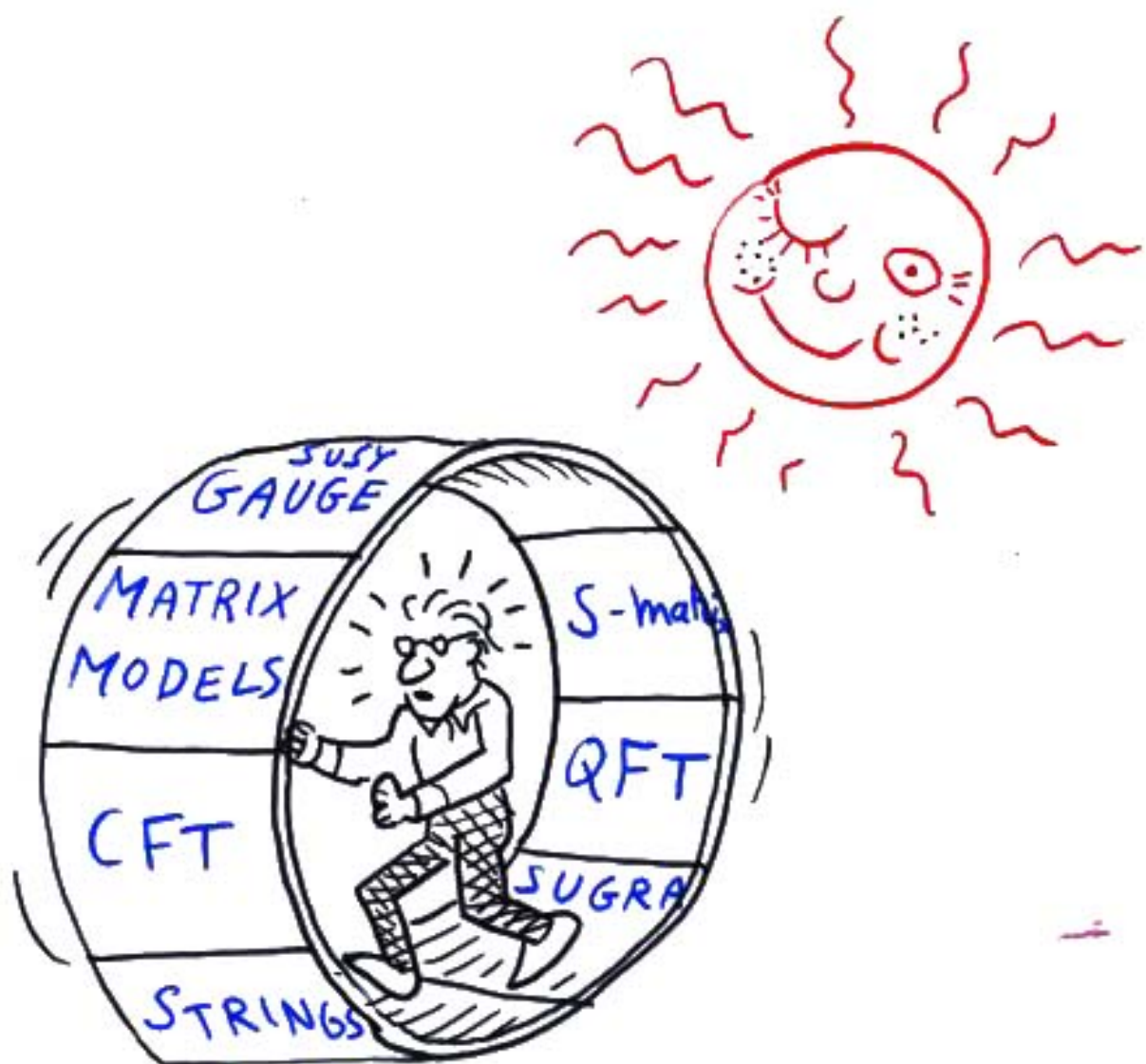
w/ Cumrun Vafa

[hep-th/0206255](#)

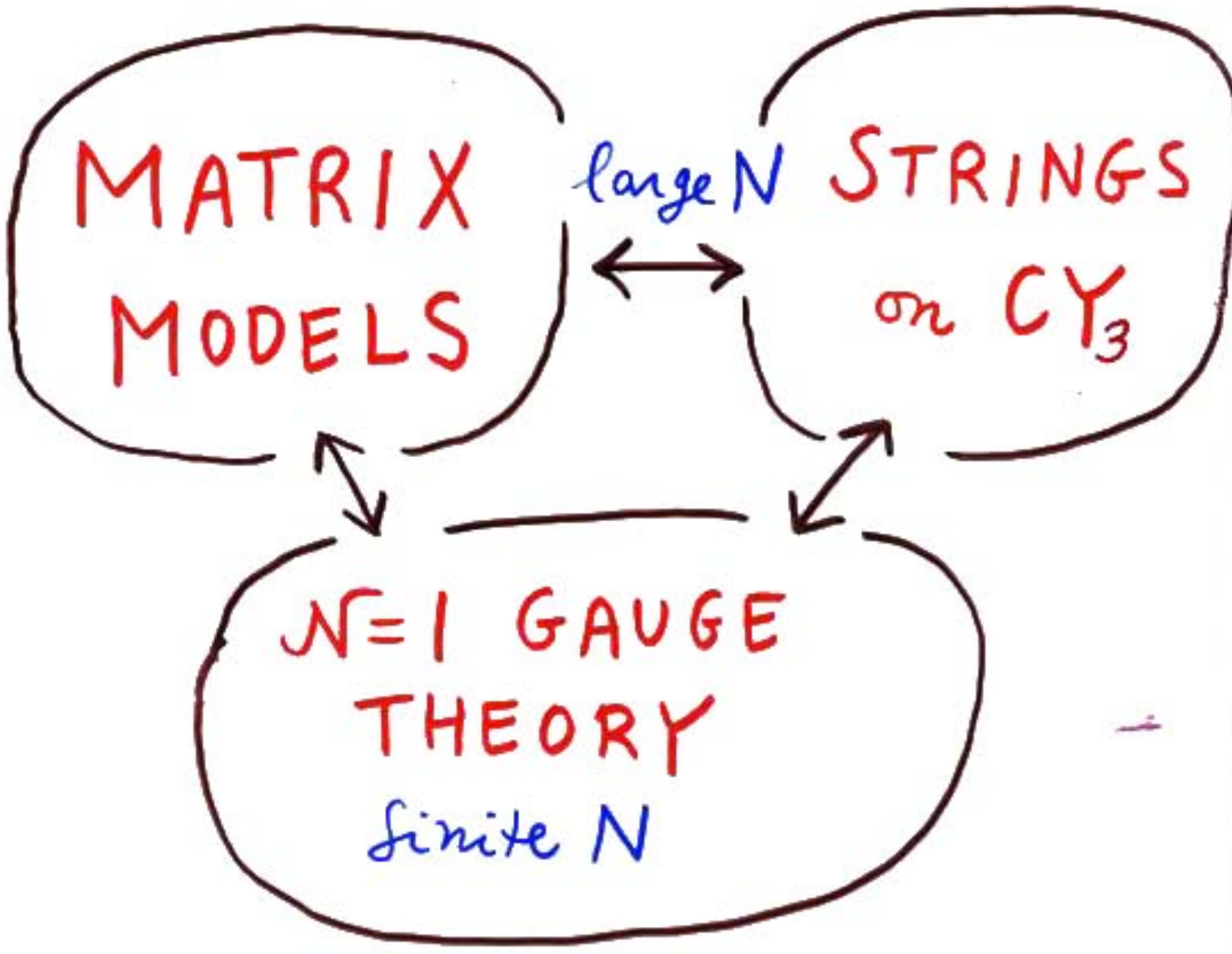
[hep-th/0207106](#)

in progress...

Vafa & collaborators  
(STRINGS 2002)



cyclic universe



# 1) Large N MATRIX MODELS

$$\int d\phi \cdot e^{\text{Tr} W(\phi)}$$

$N \times N$

~  
Saddle pt

$$e^{-\sum_{g \geq 0} N^{2-2g} \mathcal{F}_g}$$

$$\mathcal{F}_g = \sum_{\text{graphs}}$$



genus  $g$

## 2) CLOSED STRINGS in CY

$\mathcal{F}_g$ :



CY 3-fold

## 3) D=4, N=1 GAUGE THEORY

$$\int d^2\theta \text{ (F-TERMS)}$$

$$\mathcal{F}_0 \rightarrow W_{\text{eff}}(S) \text{ finite } N!$$

$$\mathcal{F}_g \rightarrow R^2\text{-terms} \text{ (grav. induced)}$$

# $N=1$ GAUGE THEORY, adj $\Phi$ <sup>5</sup>

tree level

$U(N)$

$$\int d^4x d^2\theta \cdot W(\Phi)$$

effective superpotential

$$W_{\text{eff}}(S)$$

$$S = \langle T_2 \lambda \lambda \rangle$$

gluino condensate

Burgess-Derendinger

-Quevedo-Quiroz

De la Macorra-Ross

claim:

one-loop

$$W_{\text{eff}} = N \left( S \log \left( \frac{S}{\Lambda} \right) \right) - \alpha S_{\text{bare}} + \sum_n N a_n S^n$$

matrix model

$W(\phi)$

$$\int d\phi e$$

$N \times N$

PLANAR

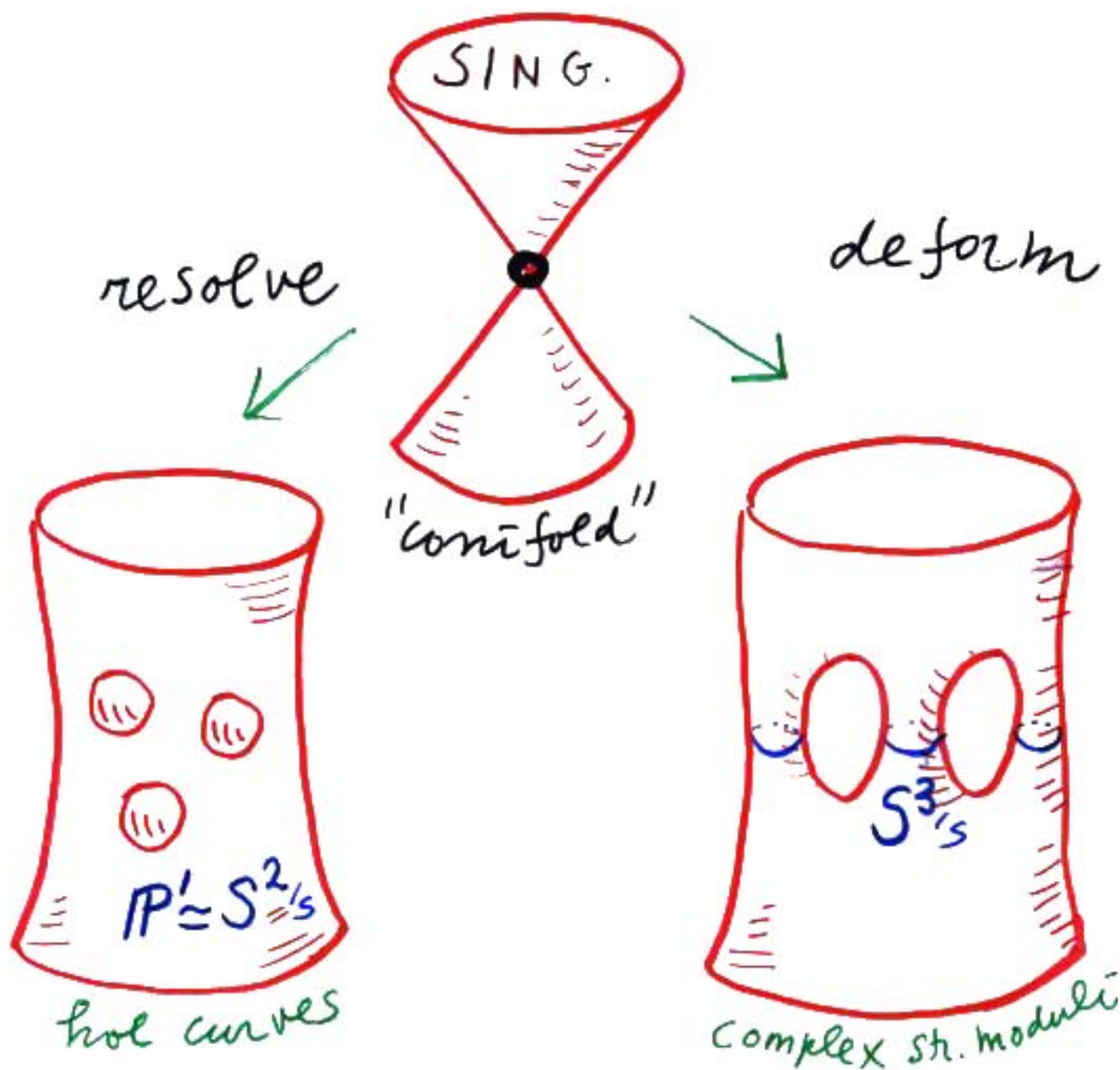
diagrams



$n$  holes

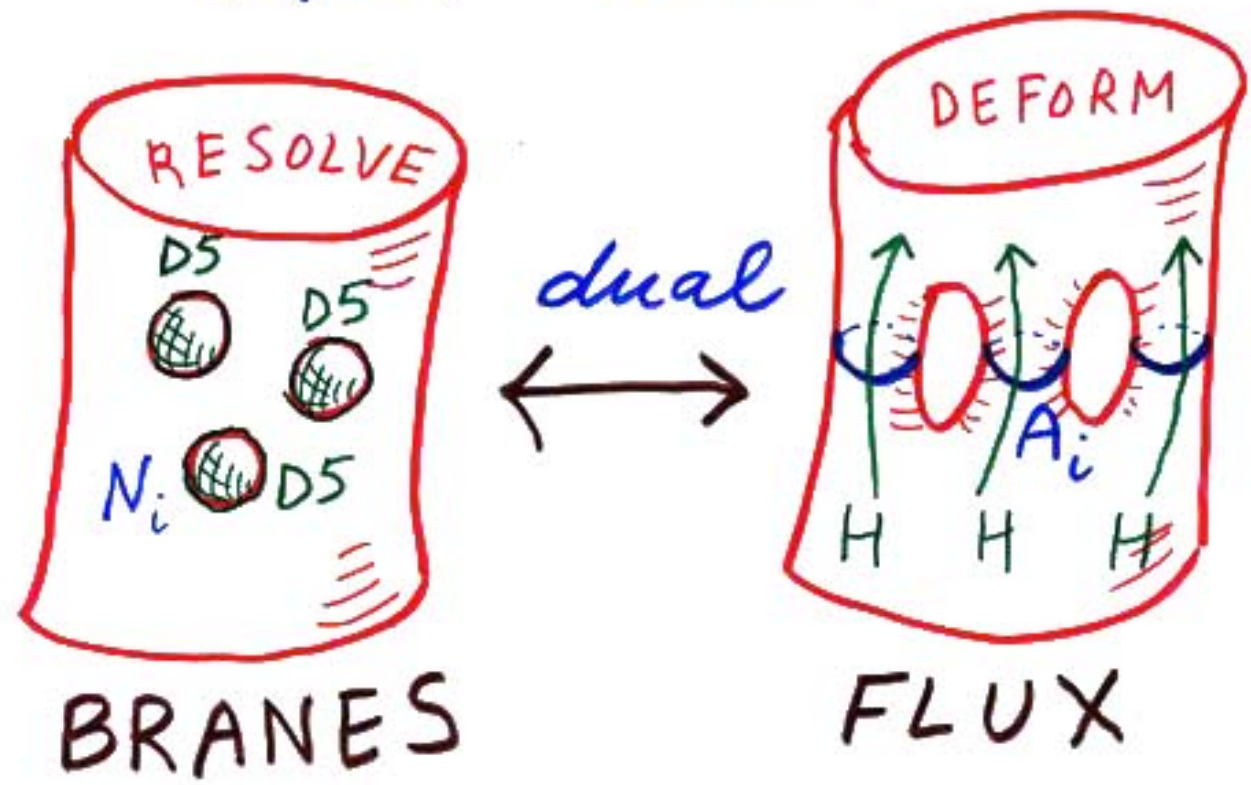
# GEOMETRIC TRANSITIONS

IIB string on  $CY_3 \times \mathbb{R}^4$



$\mathcal{N} = 2$  SUSY

# $\mathcal{N}=1$ SUSY



$$\# \text{ D5}/S_i^2 = N_i = \oint_{A_i \approx S_i^3} H_{RR}$$

$$U(N) \rightarrow U(N_1) \times \dots \times U(N_n)$$

$$\langle \text{Tr } \lambda_i \lambda_i \rangle = S_i = \text{moduli} = \oint_{A_i} \Omega$$

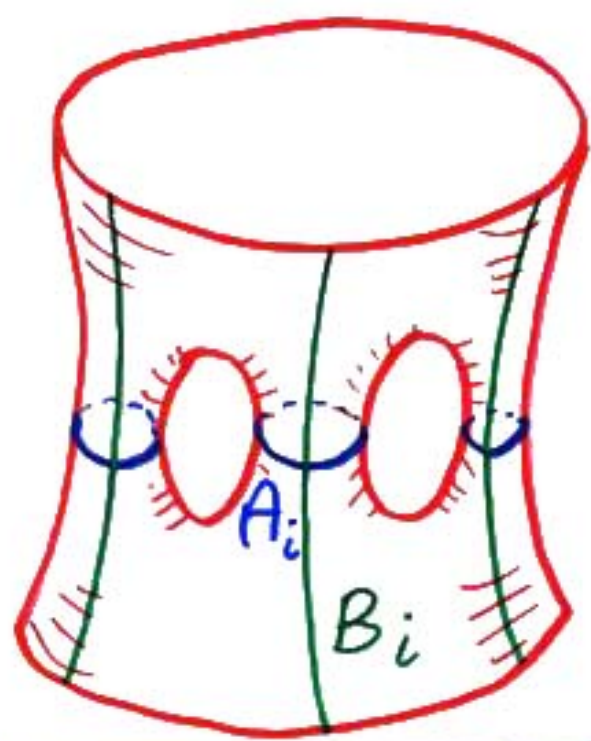
gluino condensate 3-form

# SUPER POTENTIAL

Gukov  
Vafa  
Witten

$$W_{\text{eff}}(\mathcal{S}) = \int \Omega \wedge H$$

$$= \sum_i N_i \frac{\partial \mathcal{F}_0}{\partial \mathcal{S}_i} - \alpha N_i$$



$$\mathcal{S}_i = \oint_{A_i} \Omega$$

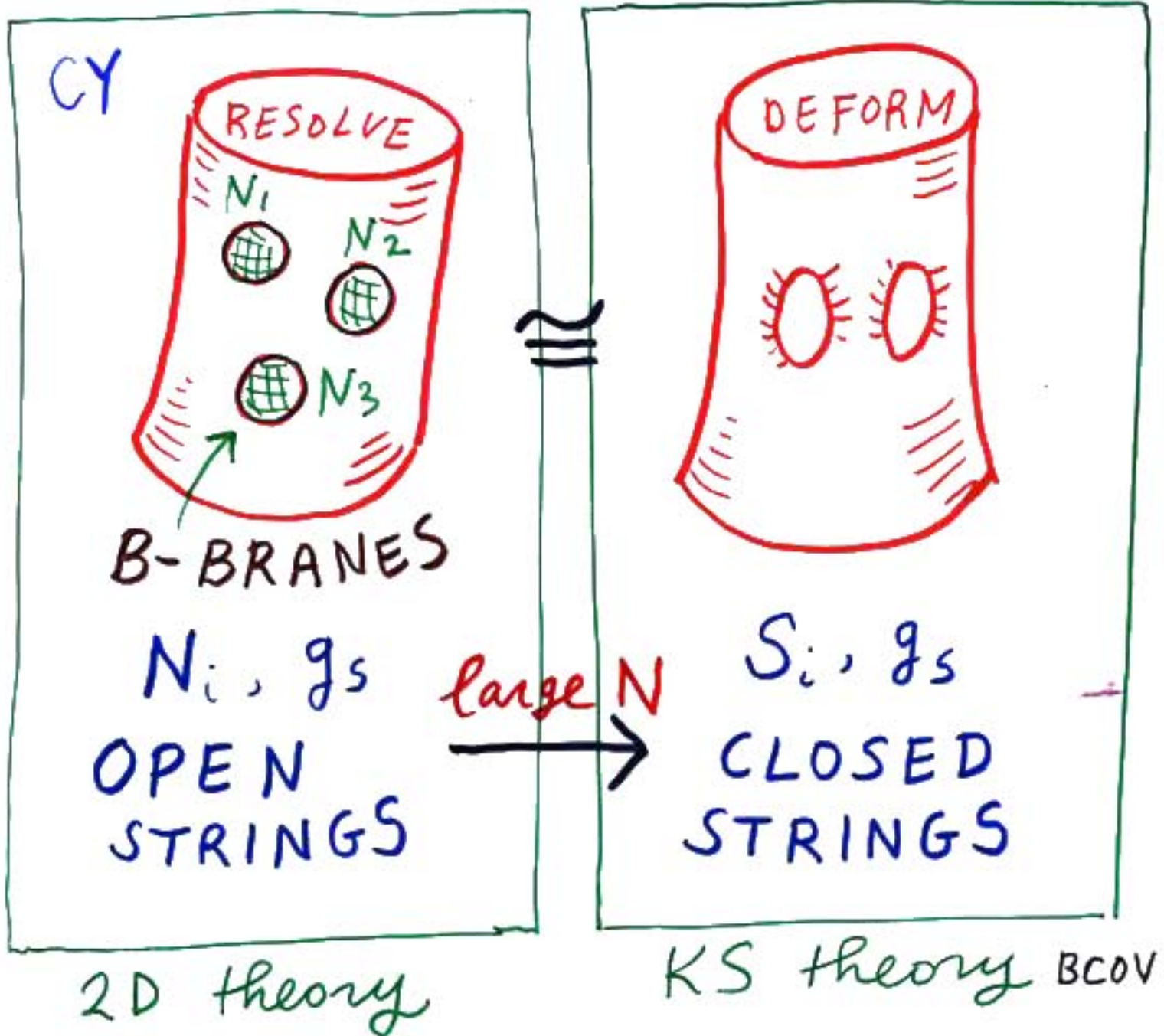
$$\frac{\partial \mathcal{F}_0}{\partial \mathcal{S}_i} = \int_{B_i} \Omega$$

prepotential  
 $\mathcal{F}_0(\mathcal{S})$

## SPECIAL GEOMETRY



## B TOPOLOGICAL STRINGS



!t Hooft couplings

$$g_s N_i = S_i = \text{moduli}$$

(mirror to A-model: Ooguri's talk)

# SIMPLE EXAMPLE

Cachazo-Intiligator-Vafa

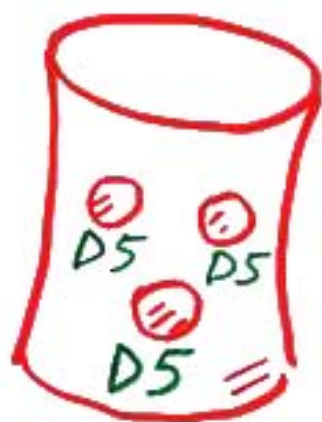
SINGULAR  $CY_3$   $W = x^{n+1} + \dots$



$$u^2 + v^2 + y^2 + W'(x)^2 = 0$$

(conifold  $u^2 + v^2 + y^2 + x^2 = 0$ )

RESOLVED :



$$\int d^2\theta \cdot T_2 W(\Phi)$$



$n$  crit pts

DEFORMED



$$u^2 + v^2 + y^2 + W'(x)^2 + f_{n-1}(x) = 0$$

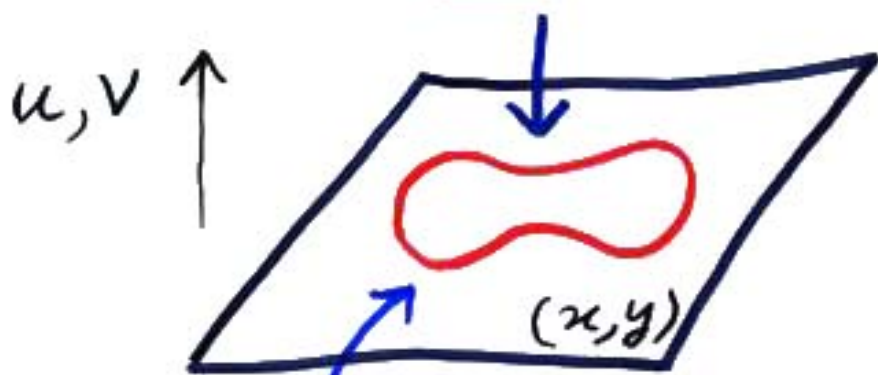
$n$  moduli  $\sim S_i$

$$(u^2 + v^2 + y^2 + x^2 = \mu)$$

# DIM<sub>al</sub> REDUCTION



CY<sub>3</sub>



complex curve

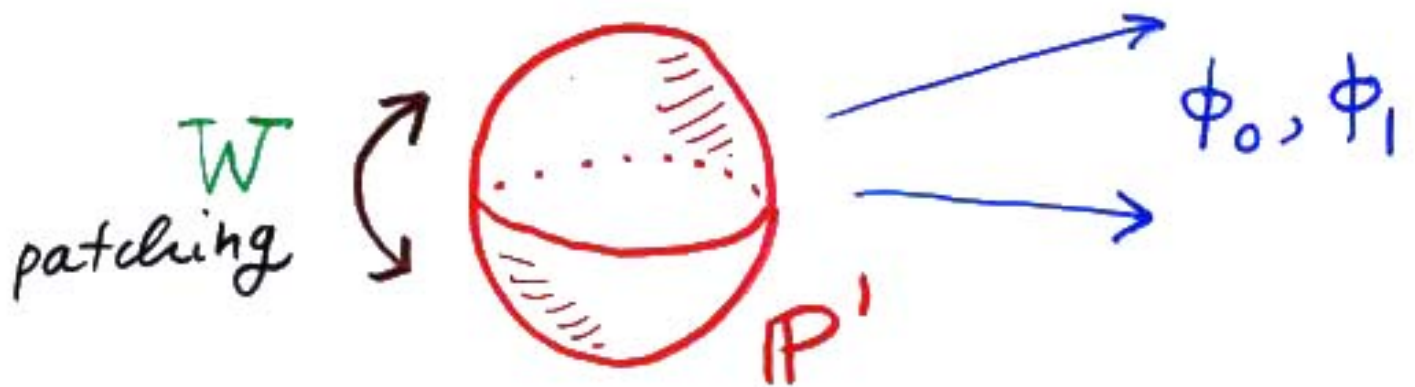
$$y^2 + W'(x)^2 + f_{n-1}(x) = 0$$

$$\Omega = \frac{du dv dx}{y} \quad \begin{array}{l} \text{3-form} \\ \text{holomorphic} \end{array}$$

$$\downarrow$$
$$y dx$$

$\begin{array}{l} \text{1-form} \\ \text{meromorphic} \end{array}$

# OPEN STRINGS on BRANE<sup>(2)</sup>



FIELDS:  $(A_{\bar{z}}, \phi_0, \phi_1)$   $N \times N$

hol CS <sup>Witten</sup>  $sp\bar{m} \uparrow_0, \uparrow_1$  - chiral

$$\int_{P^1} d^2z \operatorname{Tr} (\phi_0 \bar{D}_A \phi_1 + W(\phi_0))$$

Kachru-Katz-Lawrence  
- McGreevy

localize

$$\phi_0(z) = \phi \text{ const}$$

$$\frac{1}{\operatorname{Vol} U(N)} \int_{N \times N} d\phi \cdot e^{-\frac{1}{g_s} \operatorname{Tr} W(\phi)}$$

holomorphic

## MATRIX MODEL

# EIGENVALUE DYNAMICS <sup>13</sup>

Berezin-Itzykson-Parisi-Zuber

eigenvalues

$$\Phi = U \cdot \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_N \end{pmatrix} \cdot U^{-1}$$

$$\int \prod_i d\lambda_i \cdot \prod_{i < j} (\lambda_i - \lambda_j)^2 \cdot e^{-\frac{1}{g_s} \sum_i W(\lambda_i)}$$

JACOBIAN (integrate out  $U$ )

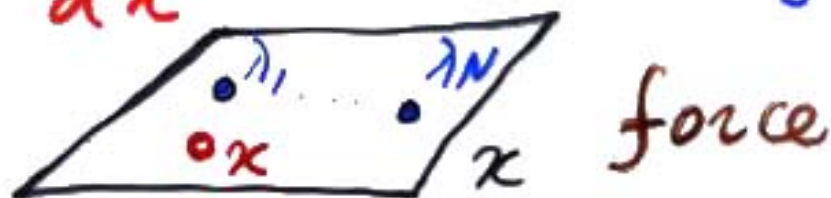
effective action

COULOMB  
repulsion

$$S = \sum_i W(\lambda_i) - 2g_s \sum_{i < j} \log(\lambda_i - \lambda_j)$$

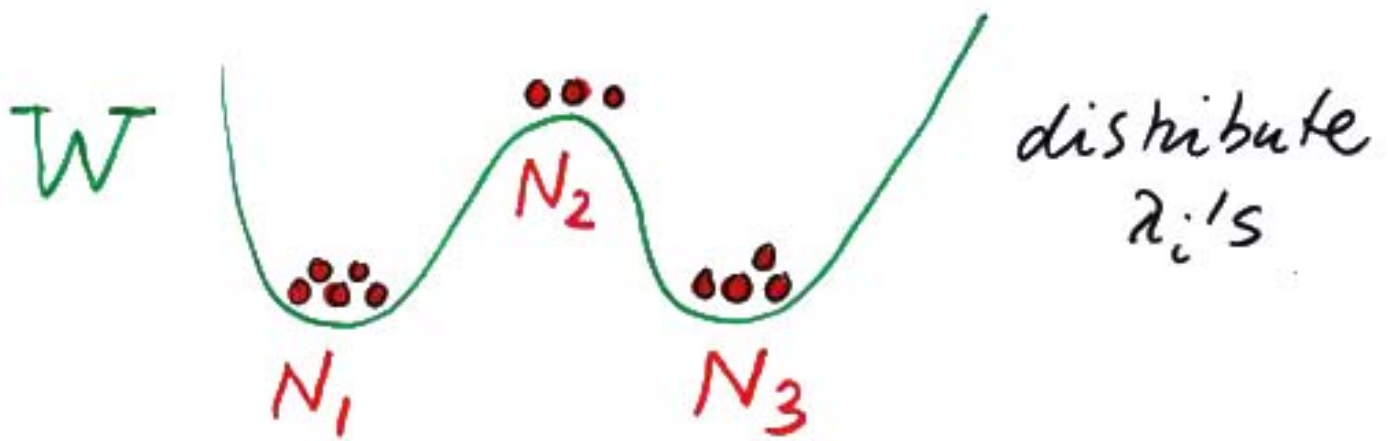
PROBE EIGENVALUE  $\lambda_i = x$

$$y(x) = \frac{dS}{dx} = W'(x) - 2g_s \sum_j \frac{1}{x - \lambda_j}$$

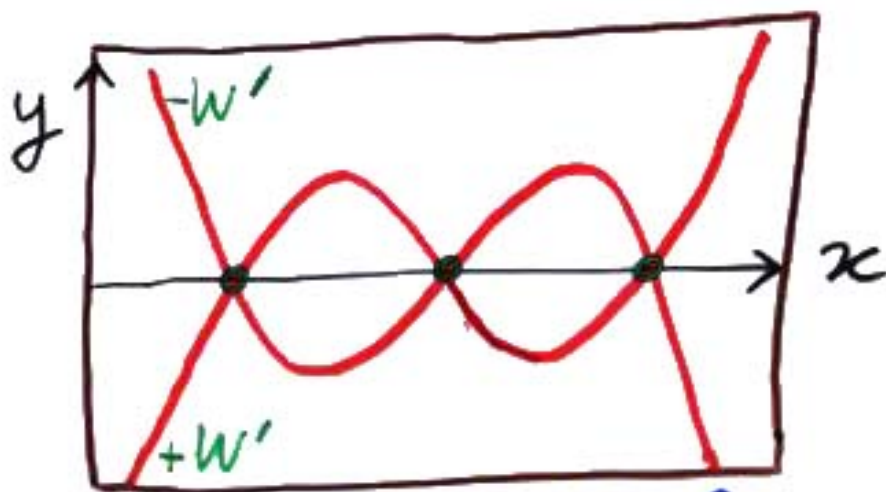
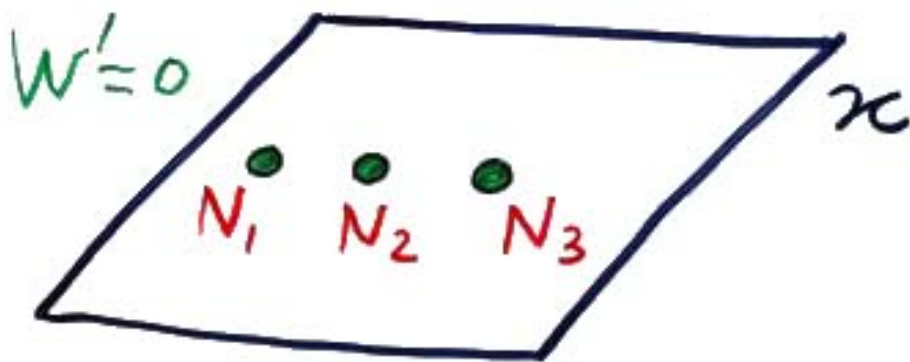


CLASSICAL  $g_s = 0$

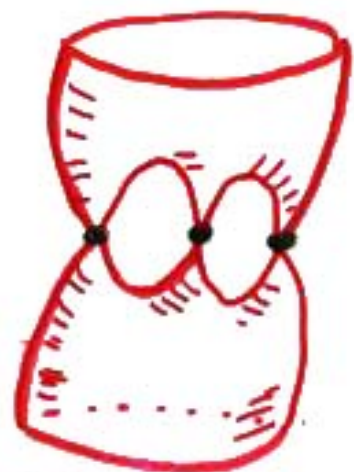
$$y = W'(x) = 0$$



filling numbers  $\sum N_i = N$



$$y^2 = W'(x)^2$$



Riemann surface

# QUANTUM $g_s > 0$

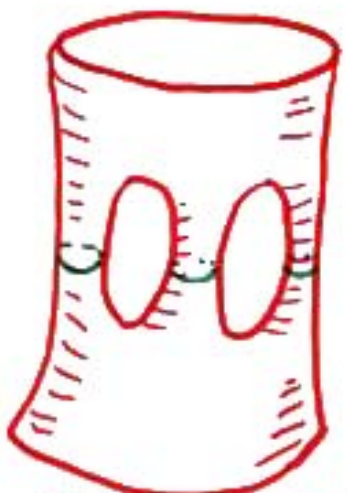
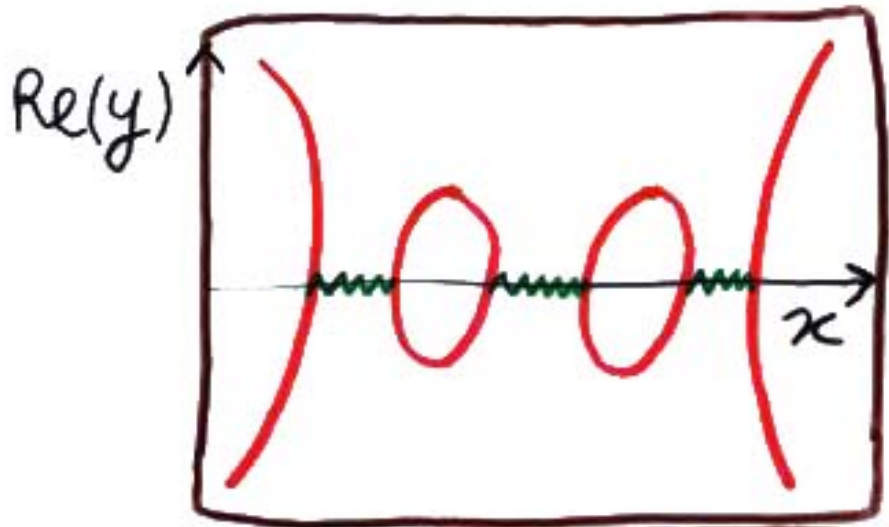
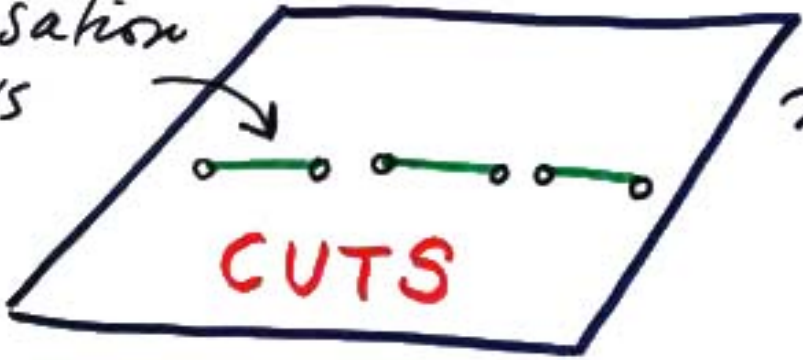
't Hooft limit  $N_i \rightarrow \infty$

$S_i = g_s N_i$  fixed  $\frac{N_i}{N}$  finite

continuous eigenvalue density

$$g(x) = \frac{1}{N} \sum \delta(x - \lambda_i) = \text{jump } \frac{y(x)}{g_s}$$

condensation of e.v.'s



Riemann surface

$y^2 = W'(x)^2 + f_{n-1}(x)$   
determined by  $S_i$

SPECTRAL CURVE

# Quantum resolution of Sing.



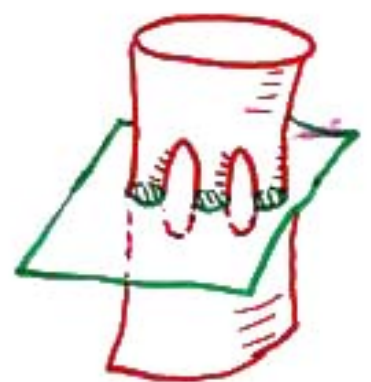
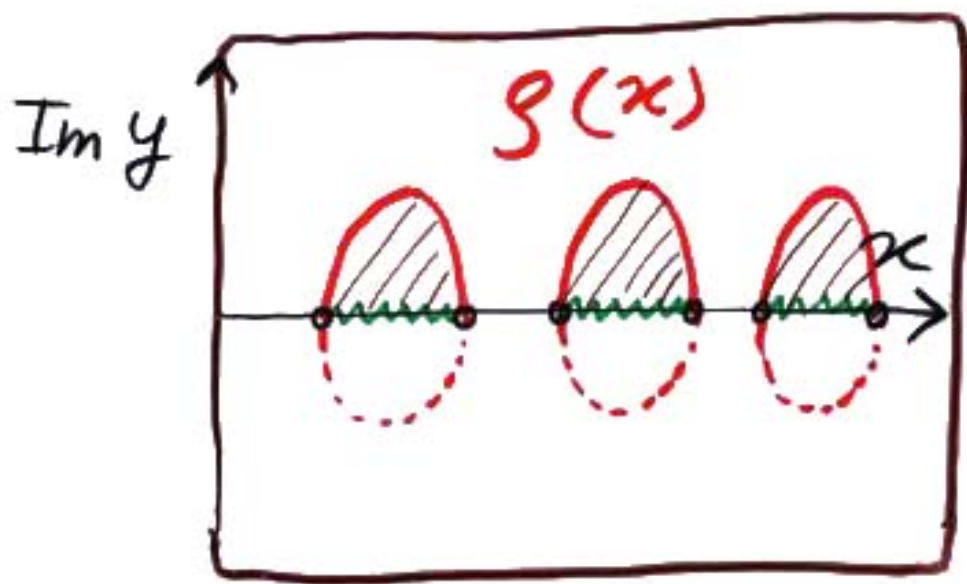
double pt



2 branch pts

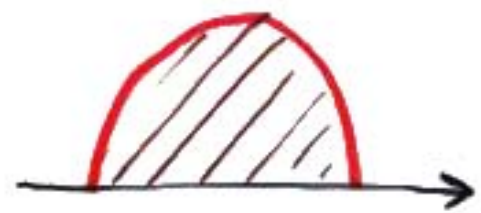
$N_i$

size cut  $\sim g_s N_i = S_i$



CONIFOLD: gaussian  $W = x^2$

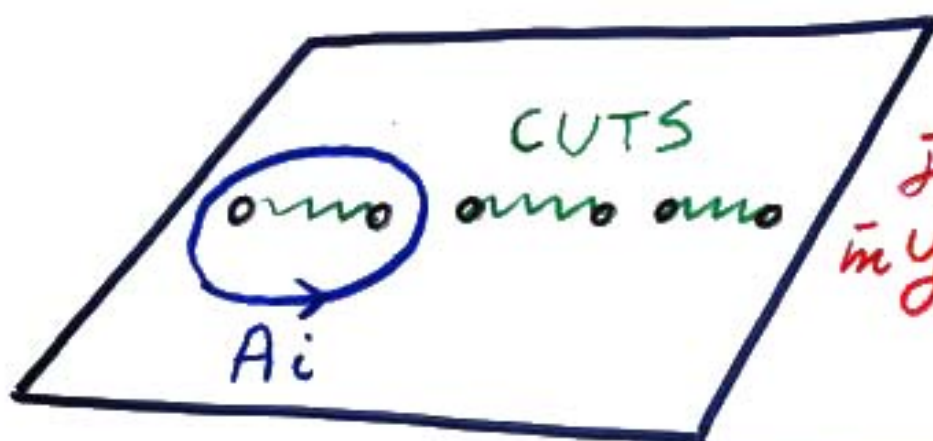
$$g(x) = \sqrt{\mu - x^2}$$



WIGNER semi-circle  $\int d\phi e^{-\tau \phi^2}$   
 $N \times N$



# SPECIAL GEOMETRY & CY<sub>3</sub><sup>17</sup>

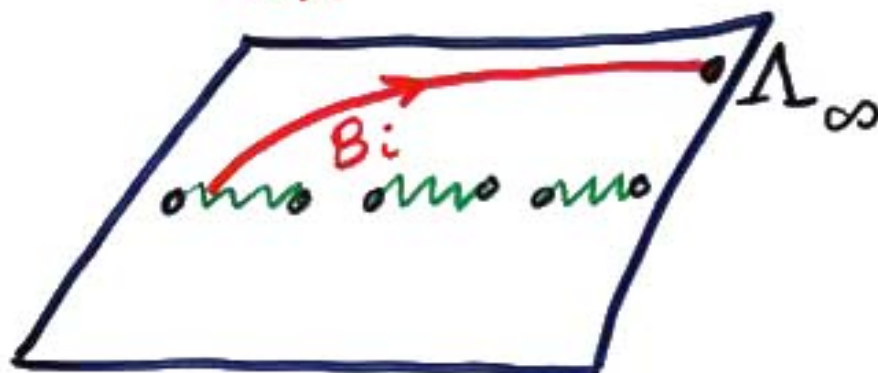


Jump  
 $\bar{m} y(x) = \int(x) g_s$

$$* \quad \frac{1}{2\pi i} \oint_{A_i} y dx = g_s N_i = S_i$$

filling #

$y(x) = \frac{dS}{dx}$  force on probe  $\lambda_i = x$



$$* \quad \int_{B_i} y dx = \frac{\partial F_0}{\partial S_i}$$

free energy

$$CY_3: u^2 + v^2 + y^2 + W'(x)^2 + f(x) = 0$$

$\Omega \rightarrow y dx$

# N=1 GAUGE THEORY

$N=2 \longrightarrow \int d^2\theta W(\phi)$



$U(N) \longrightarrow U(N_1) \times \dots \times U(N_n)$

$\Downarrow$  confine.

gluino condensates

$S_1, \dots, S_n$

## EFFECTIVE SUPERPOTENTIAL

CIV

$$W_{\text{eff}} = \sum_i \left( N_i \frac{\partial \mathcal{F}_0}{\partial S_i} - \alpha S_i \right)$$

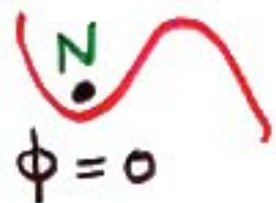
*vol  $U(N_i)$  in MM*

$= \sum N_i S_i \log S_i / \Lambda +$  PERTURBATIVE

*"one-loop"*

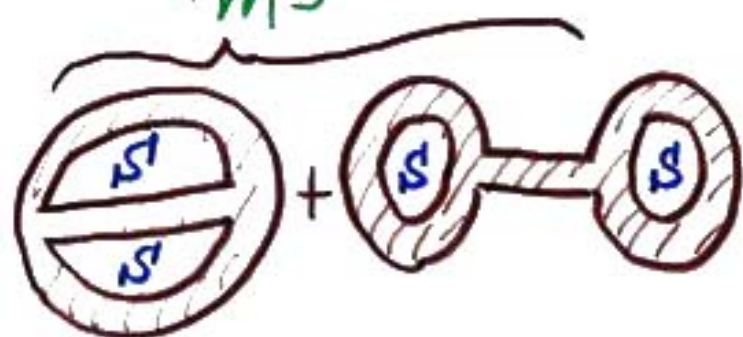
only PLANAR diagrams contribute!

e.g.  $W = m\phi^2 + g\phi^3$ ,



$W_{eff} = N S \log S/\Lambda - \alpha S$   
↑ bare coupling

$+ * \frac{g^2}{m^3} S^2 + \dots$  CIV



← computed by MATRIX MODEL

SOLVE:  $dW_{eff} = 0$

$S \sim e^{-\frac{\alpha}{N}\Lambda}$   
+ exact inst.

$W_{eff} = \sum$  instantons

- PERTURB.  $\Rightarrow$  INSTANTONS  
nth order  $\sim$  instantons (Nekrasov?)
- PLANAR = EXACT
- Large N = CY 3-fold  
"master field"

CLAIM: this works in general\*

\* classical groups, vectorlike rep, massive + abelian vacua.

D=4, N=1 gauge theory

$$\int d^2\theta \cdot W(\phi_i)$$

MATRIX MODEL:  $\Rightarrow$  0 dim

$$\int d\phi_i \cdot e^{-\frac{1}{g^2} W(\phi_i)}$$

PLANAR:  $W_{eff}(S) = \sum (N_i \cdot \alpha_i \cdot F_0 - \alpha_i S_i)$

NON-PLANAR: gravity induced

e.g.  $\mathcal{F}_1(S) \cdot R \wedge R$  

WHY?

field theory argument

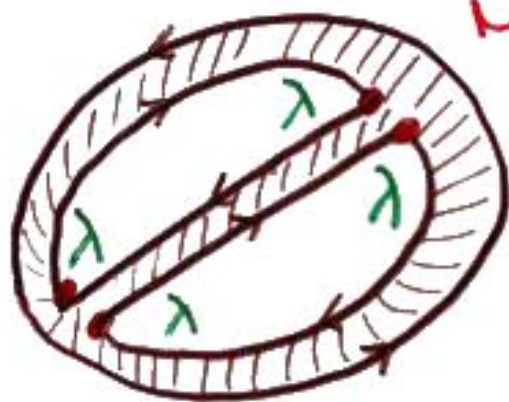
Sketch (in progress)

$$\int d^2\theta \cdot W_{\text{eff}}(\mathcal{S}) \quad \text{"topological" INDEX}$$

F-term: localize  $x^\mu = \theta^\alpha \text{const}$

$$\int d^4x d^2\theta W_{\text{tree}}(\Phi) \xrightarrow{\text{QFT}_4} W(\Phi) \text{ MM}_0$$

PLANAR?  $\lambda\lambda$  / boundary (cut)



top strings

BCOV  
Berkovits

$$\alpha' \rightarrow 0$$

Schwinger picture

QFT

# QUIVERS: $ADE, \hat{A} \hat{D} \hat{E}$

$$W = \sum_i W_i(\phi_i) + \sum_{\langle i, j \rangle} Q_{ij} \phi_j Q_{ji}$$

Cachazo-Katz-Vafa

Solvable multimatrix models

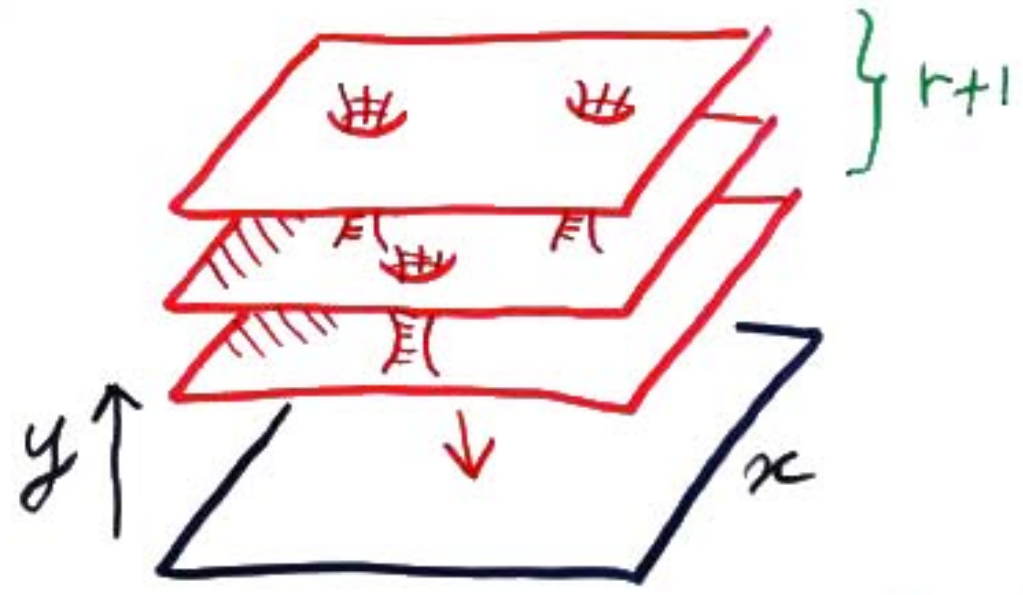
Kostov  
Kharchev et al

strings on



$$\Pi(\lambda_I - \lambda_J)^2 \rightarrow \Pi(\lambda_{i, I} - \lambda_{j, J})^{P_i \cdot P_j}$$

## MULTIPLE COVERS



$$y^{r+1} + x^n + \dots + u^2 + v^2 = 0$$

# SEIBERG-like dualities

(same univ. class MM)

(in progress)

MM  $\longrightarrow$  CY<sub>3</sub> > curves

e.g.  $W = \phi_1^n + \phi_1 \phi_2^2$

$\downarrow$  CKV

$$u^2 + F(v, x, y) = 0$$

complex surface  
("Laufer geometry")



$$\underline{N=4 \rightarrow N=1^*} \quad (\text{in progress})$$

$$W = \phi_1 [\phi_2, \phi_3] + m \sum \phi_i^2$$

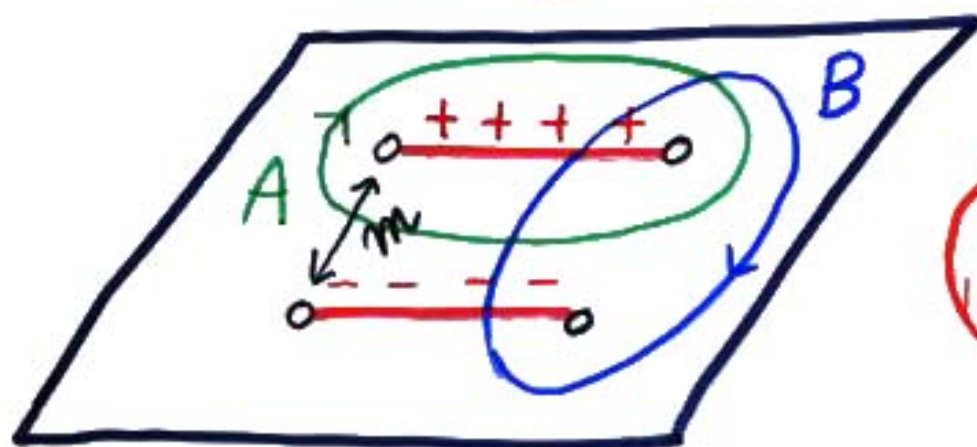
3 MATRIX MODEL

Kazakov  
- Kostov  
- Nekrasov

integrate out  $\phi_2, \phi_3$

$$\int d\phi_1 \frac{e^{-T_2 \phi_1^2}}{\det(m + [\phi_1, -])}$$

↓ large N



elliptic curve

S<sup>1</sup>-DUALITY

periods

$$\Pi_A = \oint_A y dx$$

$$\Pi_B = \oint_B y dx$$

$$W_{\text{eff}} = N \Pi_B - \tau_0 \Pi_A$$

↓ minimize

$$\tau = \frac{\partial \Pi_B}{\partial \Pi_A} = \frac{\tau_0}{N}$$

$$W_{\text{eff}} \sim E_2(\tau) + \dots \quad \text{Dorey}$$

$SL(2, \mathbb{Z})$  perturbative

# LEIGH-STRASSLER deformation <sup>28</sup>

$$[\phi_2, \phi_3] \rightarrow \phi_2 \phi_3 - q \phi_3 \phi_2$$

$$q = e^{2\pi i \beta}$$

## 6 VERTEX MODEL

Kazakov  
- Kostov

$$\beta \rightarrow \beta + \tau$$



preserve modularity

$$\beta \rightarrow \beta / (c\tau + d)$$

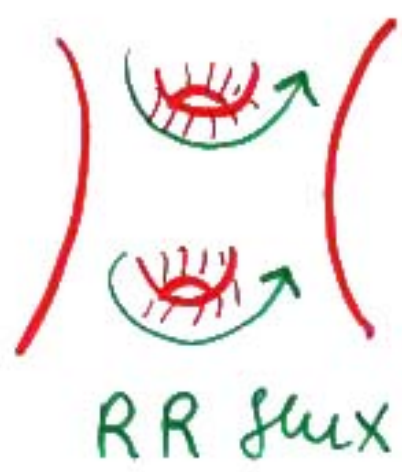
# DOUBLE SCALING LIMITS

$g_s$  fixed,  $N \rightarrow \infty$

e.g.  $N_1 - N_2$  fixed (cf  $PP \approx$ )

## "OLD" (1989) MATRIX MODELS

$y^2 + x^3 + \dots = 0$  pure gravity



## UNITARY MODELS $\Phi \approx \Phi + 2\pi$ periodic

$$U = e^{i\Phi}, \quad W = \cos \Phi$$

GROSS-WITTEN  $\rightarrow$  SEIBERG-WITTEN

$$\int dU \cdot e^{\text{Tr} U + \bar{U}}$$

one-plaquette

$$\mathcal{N} = 2 \text{ SU}(2)$$

# CONCLUSION:

$$\int_{N \times N} d\phi \cdot e^{T_2 W(\phi)} \sim e^{\sum N^{2-2g} F_g}$$

↙  
CLOSED  
STRINGS



CY<sub>3</sub>

# HIGHER GENUS

non-linear LOOP EQNS

$$y = \partial\varphi(x)$$

$$\langle \partial\varphi(x)^2 \rangle = W'(x)^2 + f(x)$$

↓ quantum CY

$\langle T(x) \rangle$  regular

VIRASORO constraints

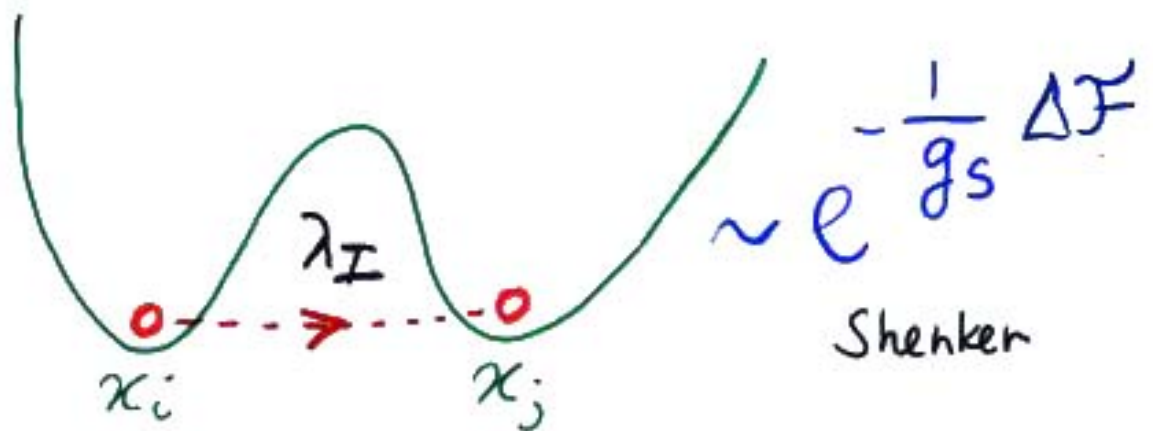
$$L_n Z = 0 \quad n \geq -1$$

reparametrize matrix  $\Phi$

# EIGENVALUE TUNNELING

$\cong$

## DOMAIN WALLS



$$\text{tension} \sim \frac{1}{g_s} (\partial_{x_i} \mathcal{F} - \partial_{x_j} \mathcal{F}) \sim \Delta \bar{W}$$

