

On HS symmetries, cubic interactions, AdS & CFT

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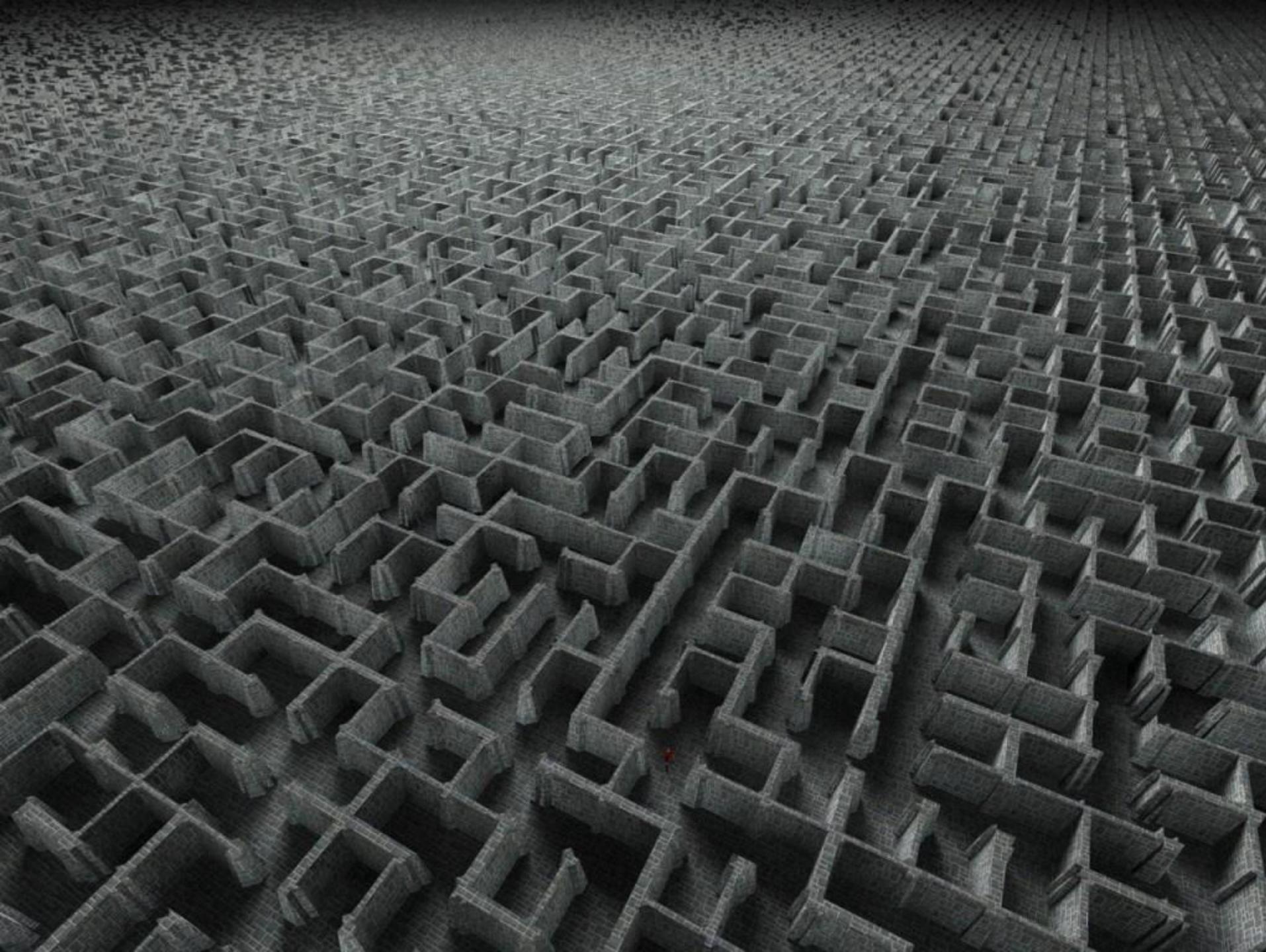
MAX-PLANCK-GESELLSCHAFT

based on arXiv:1311.0242 w. E.Joung and arXiv:1305.5180 w. N.Boulanger, D.Ponomarev, E.Skvortsov

Desired Goals!



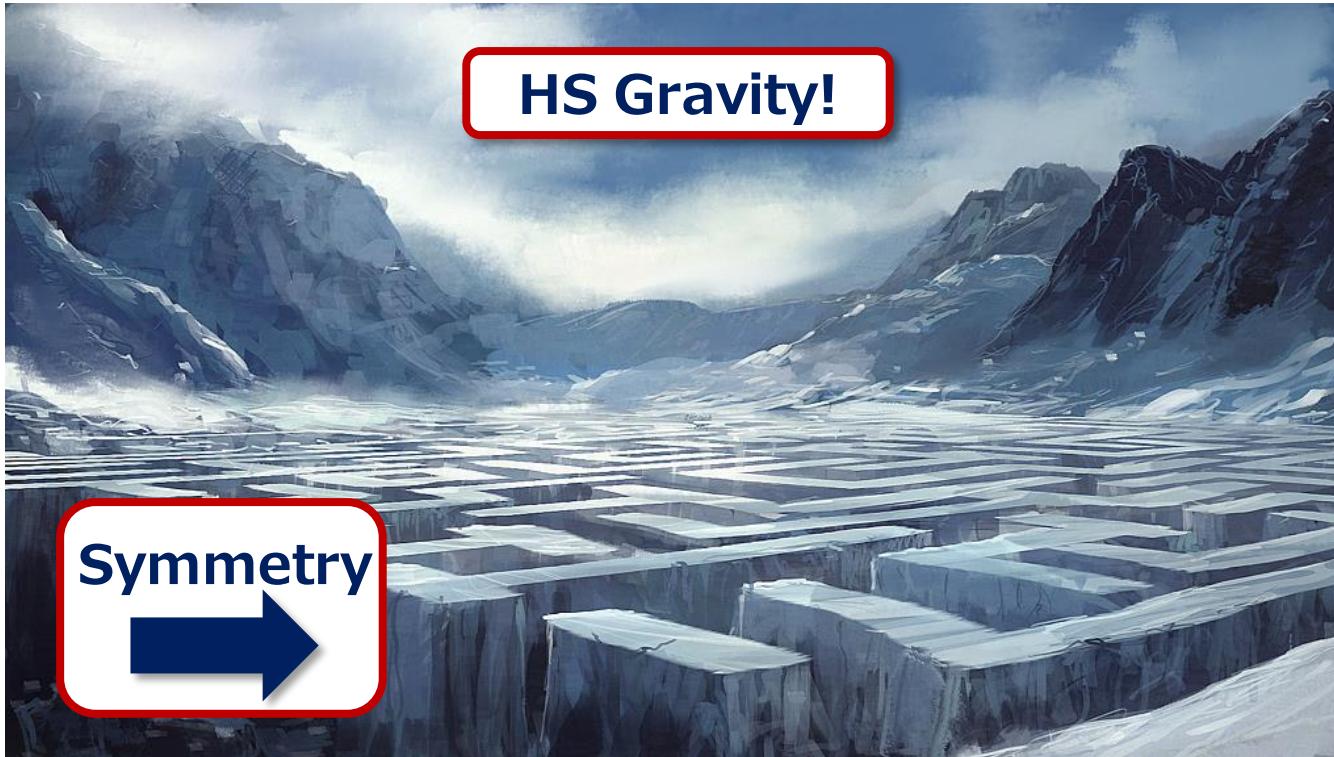
- **Uncover the symmetries behind String Theory**
- **Shed light on holography and Quantum Gravity**



...topology matters!!



History teaches us: Symmetry has always paid off!



What is the “maximal” symmetry of Quantum Gravity?

Fradkin & Vasiliev ('80s): Higher-Spin Symmetry!

Cons:

too big symmetry => too simple theory...

Pros:

- Deformation (gauging) of symmetries (Algebroid, field dependent structure constants)
- Almost all checks of AdS/CFT driven by symmetry (maximal symmetry => proof of Holography?)

To do List

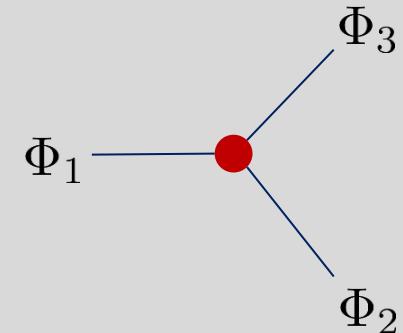
- HS systematics
 - HS holography
 - HS on non-trivial backgrounds
(black-holes!)
 - ...
- Some of this today



Cubic Couplings Classification

$$s_1 \geq s_2 \geq s_3$$

	n	$\#\partial$	$\delta_{E_1}^{(1)}$	$\delta_{E_2}^{(1)}$	$\delta_{E_3}^{(1)}$	$C^{(3)}$
Class I	0	$s_1 + s_2 + s_3$	= 0	= 0	= 0	$\approx \tilde{K}(Y_\ell, H_{12}, H_{23}, H_{31})$ $\ell = 2 \text{ or } 3$
	\vdots	\vdots	\vdots	\vdots	\vdots	
	$\frac{s_2+s_3-s_1}{2}$	$2s_1$	= 0	\vdots	\vdots	
Class II	\vdots	\vdots	$\neq 0$	\vdots	\vdots	$\approx \tilde{K}(Y_1, H_{12}, H_{23}, H_{31})$
	\vdots	\vdots	\vdots	\vdots	\vdots	
	$\frac{s_3+s_1-s_2}{2}$	$2s_2$	\vdots	= 0	= 0	
Class III	\vdots	\vdots	\vdots	$\neq 0$	Λ	
	\vdots	\vdots	\vdots	\vdots	\vdots	
	$\frac{s_1+s_2-s_3}{2}$	$2s_3$	\vdots	\vdots	Λ	
Class IV	\vdots	\vdots	\vdots	\vdots	$\neq 0$	
	s_3	$s_1 + s_2 - s_3$	$\neq 0$	$\neq 0$	$\neq 0$	



$$[E_1, E_2]_3^{(0)} = 0 \iff \delta_{E_1}^{(1)} = 0 \text{ or } \delta_{E_2}^{(1)} = 0$$

E.Joung and M.T. '13

Coleman-Mandula in AdS/CFT!

Assumptions: symmetric tensors + one HS generator + no colour (+ Gravity)!

Boulanger, Ponomarev, Skvortsov, M.T. '13

AdS₃/CFT₂
(Virasoro!!!)

AdS₄/CFT₃

AdS₅/CFT₄

AdS_{6+n}/CFT_{5+n}

Algebra:

$$\text{hs}_3(\lambda) = \frac{U[sl(2)]}{C_2 - \lambda 1}$$

Moyal (unique)

$$\text{hs}_5(\lambda) = \frac{U[su(2, 2)]}{C_2 - \lambda 1}$$

$$\text{hs}_d = \frac{U[so(d-1, 2)]}{\square \square \oplus \square \square \square}$$

UIR

$$(\square - m_\lambda^2)\phi(x) = 0$$

**Scalar & Spinor
Singletons**

Fradkin, Vasiliev '86; Konstein, Vasiliev '90; Maldacena, Zhiboedov 2011

**One parameter family
(doubletons)**

Gunaydin et all; Boulanger, Skvortsov; Manvelyan, Mkrtcyan, Mkrtcyan, Theisen

Scalar singleton

$$\langle J \cdots J \rangle = \text{Tr}[\psi \star \cdots \star \psi]$$

Outlook

- Admissibility will put further constraints
- Deformation (Lie Algebroids!)
- More general backgrounds (Black-Holes!)



