

Bulk causal features + boundary correlation in AdS/CFT

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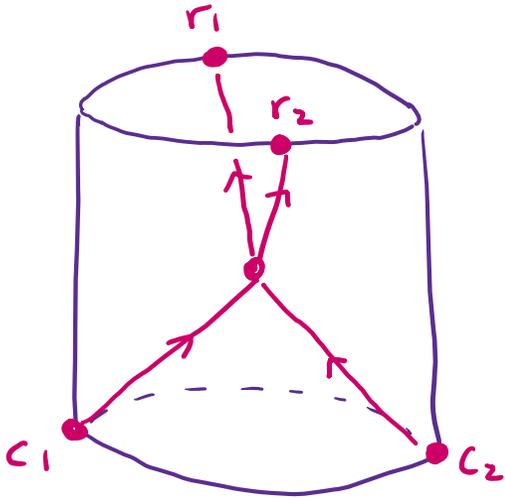
The University of British Columbia

Based on: 2105.08094 AM
2101.08855 AM

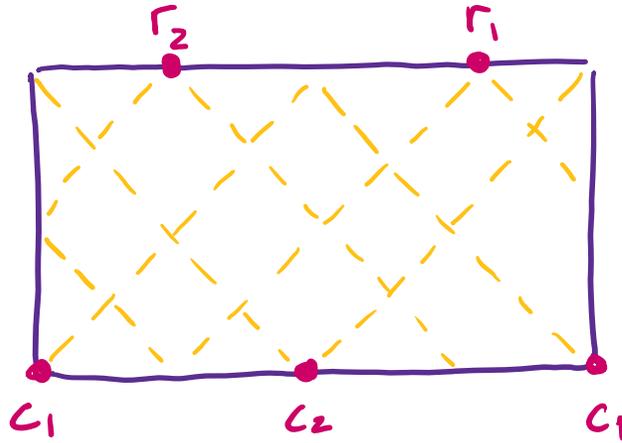
See also: 1912.05649 AM, Penington,
Sorce
1902.06845 AM

Bulk causal features

- A classic, fruitful, set-up in AdS/CFT:



can "scatter" in bulk



no place to scatter

- HPS understood that for the boundary to reproduce bulk physics, boundary theory must be large N , gapped

This talk:

what do bulk causal features tell us about the boundary state?

This talk

How do bulk causal features
constrain the boundary state?

- ① Recall ways to measure correlation and relationship (holographically) to extremal surfaces
 - ② State the "Privacy-duality" theorem and give argument
 - ③ State connected wedge theorem
 - ④ Final comments
- } Follow from dynamics of quantum information in spacetime

Measuring correlations in quantum states

Q.I. Measures of correlation

- We will be concerned with:

① Mutual information: $I(A:B) \equiv S(A) + S(B) - S(AB)$

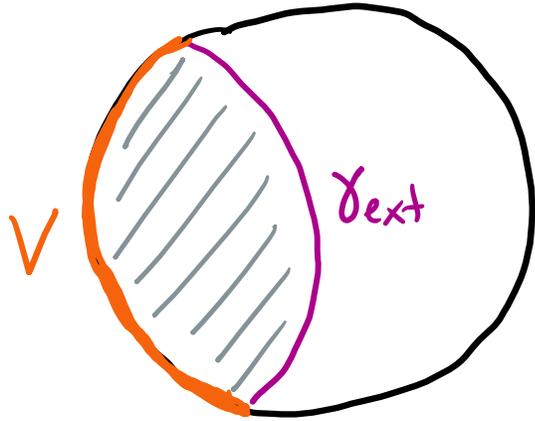
↳ Amount of correlation between A and B

② Conditional mutual info: $I(A:C|B) \equiv I(A:BC) - I(A:B)$

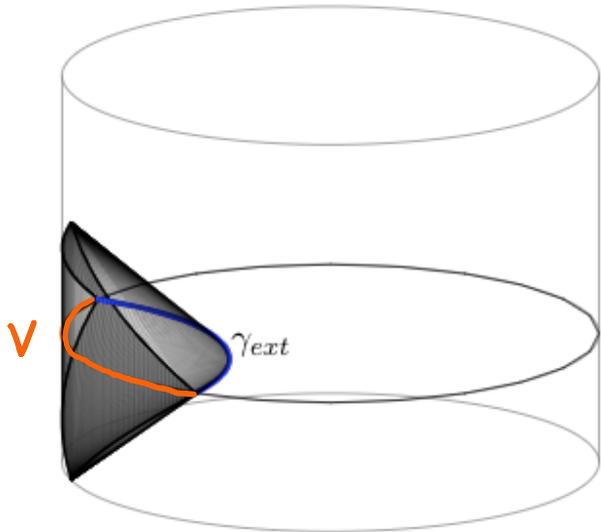
↳ Interpret as amount of correlation between
between A and BC not already present between A and B

RT Formula + the entanglement wedge

- Given a boundary region V :



$$- S(V) = \frac{\text{Area}[\gamma_{\text{ext}}]}{4G_N} + S[E_V]$$



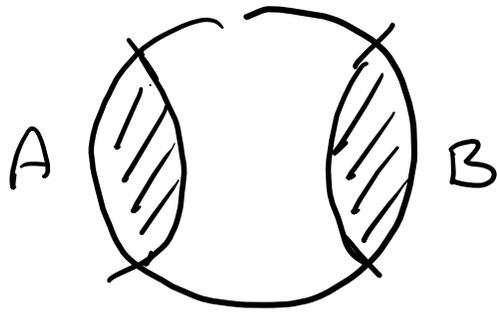
- \mathcal{P}_V records all the bulk data inside the entanglement wedge E_V

↳ "entanglement wedge reconstruction"

Correlation + extremal surfaces

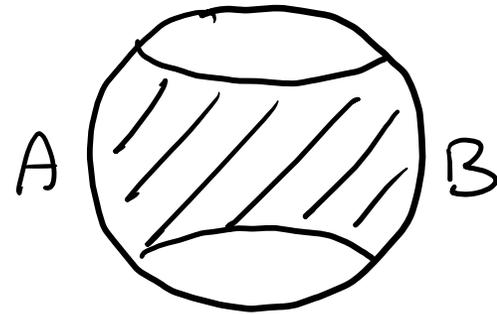
- The order of these correlation measures records qualitative features of extremal surfaces

E.g. $I(A:B) = O(1)$



$$\bar{E}_{AB} = E_A \cup E_B$$

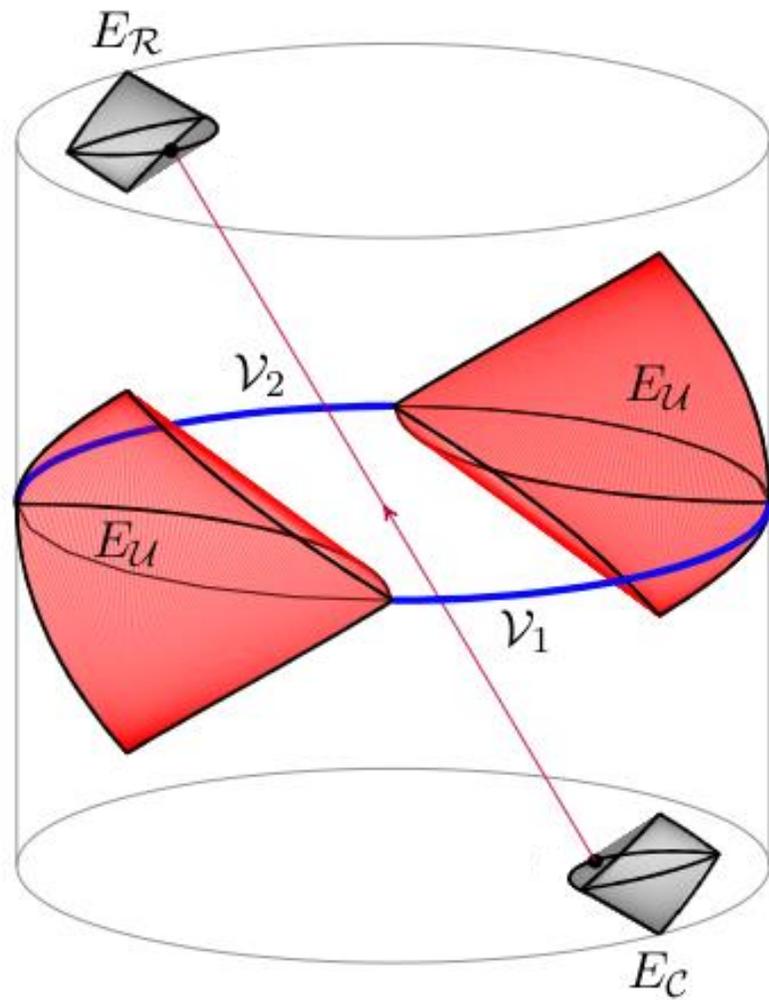
$$I(A:B) = O(1/G_N)$$



$$E_{AB} \neq E_A \cup E_B$$

Privacy-duality theorem

Privacy - duality



"private" curve

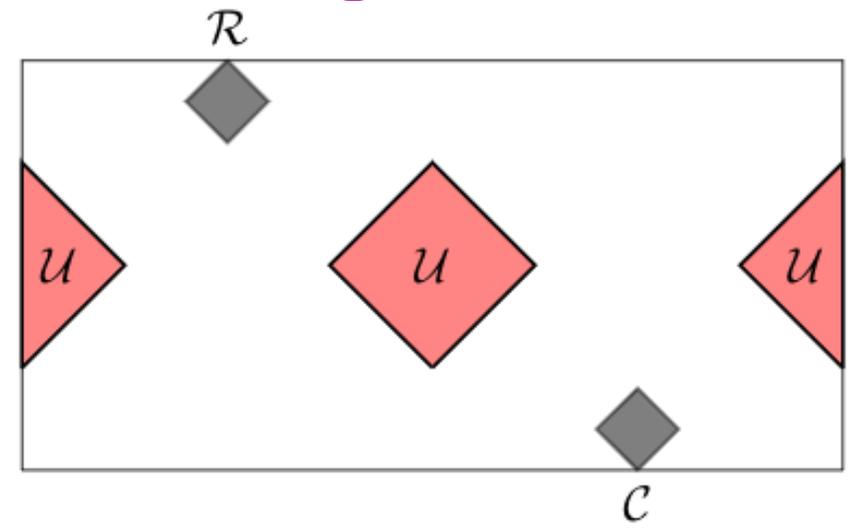


If there's a causal curve from E_c to E_R that avoids E_U ,
then

$$I(V_1:V_2|U) = O(1/GM)$$

Privacy - Duality theorem

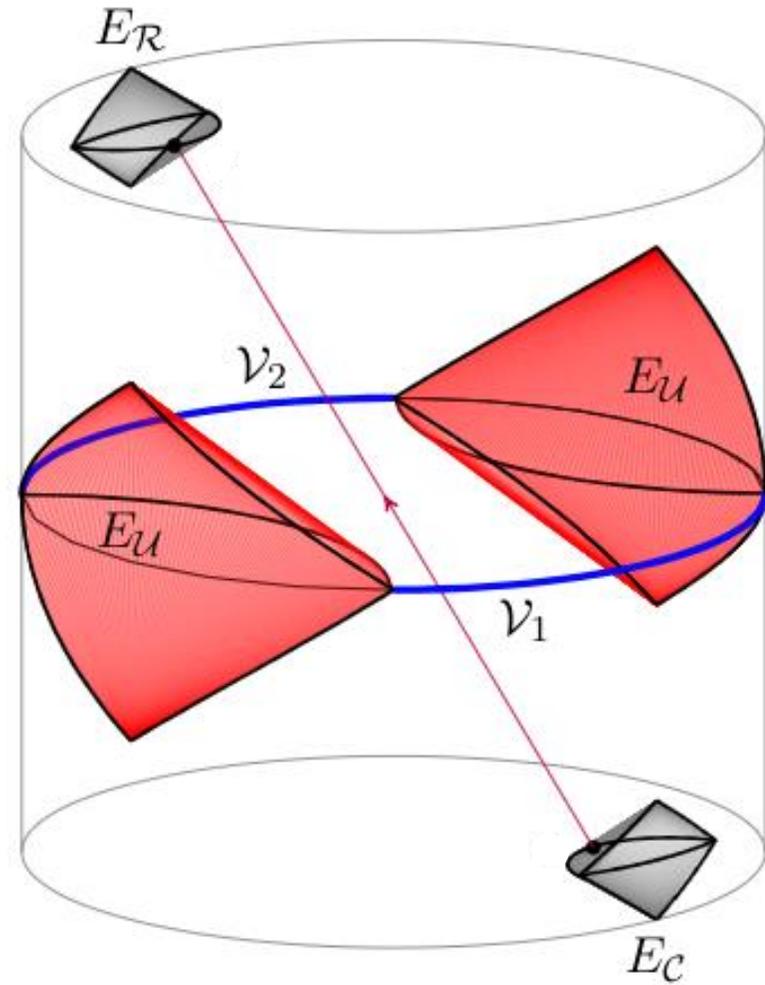
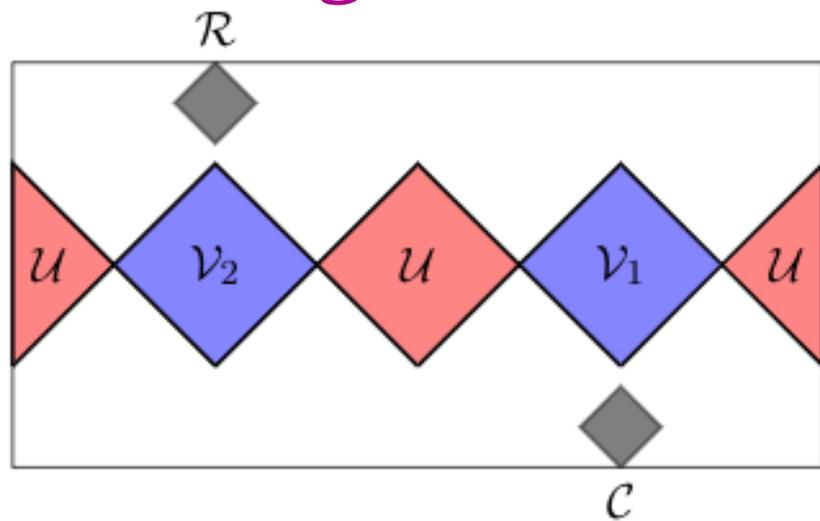
- Pick e, R, U in boundary, and a Cauchy surface Σ "through" U



Define: $V_1 \equiv \hat{D}(\mathcal{I}^+(e) \cap \Sigma \cap U')$
 $V_2 \equiv \hat{D}(\mathcal{I}^-(R) \cap \Sigma \cap U')$

Privacy-Duality theorem

- Pick e, R, U in boundary, and a Cauchy surface Σ "through" U



Define: $V_1 \equiv \hat{D}(\mathcal{J}^+(e) \cap \Sigma \cap U')$

$V_2 \equiv \hat{D}(\mathcal{J}^-(R) \cap \Sigma \cap U')$

Then:

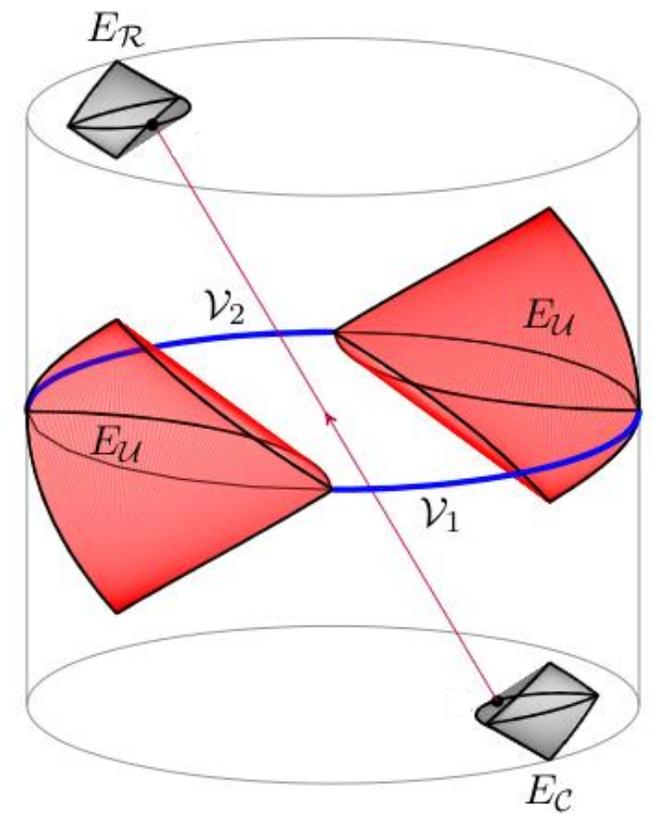
IF \exists a causal curve from E_e to E_R that avoids E_U , then $I(V_1:V_2|U) = \alpha(1/\delta n)$

Proof strategy (general)

- ① Introduce a probe system Q (or systems Q_i) to the bulk, and exploit causal feature to have Q evolve in some interesting way
- ② Translate evolution of Q to a boundary statement
- ③ Argue that for boundary to realize this evolution, while lacking same causal feature, appropriate correlations are necessary

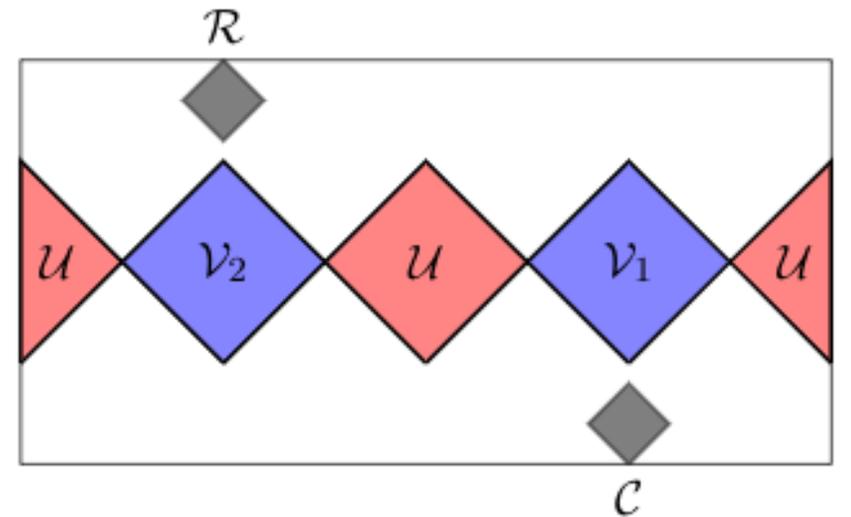
Q.I. argument for Privacy-duality

- ① Introduce a "probe" system Q , holding state $|\psi\rangle_Q$, and send it along Π
- ② Sends message from E_c to E_R , while keeping message secret from E_U
 - EW reconstruction \rightarrow Sends message from E to R , while keeping message secret from U



In boundary, there is no private curve.

- ③ \hookrightarrow Instead, exploit correlations among subsystems to hide state on Q



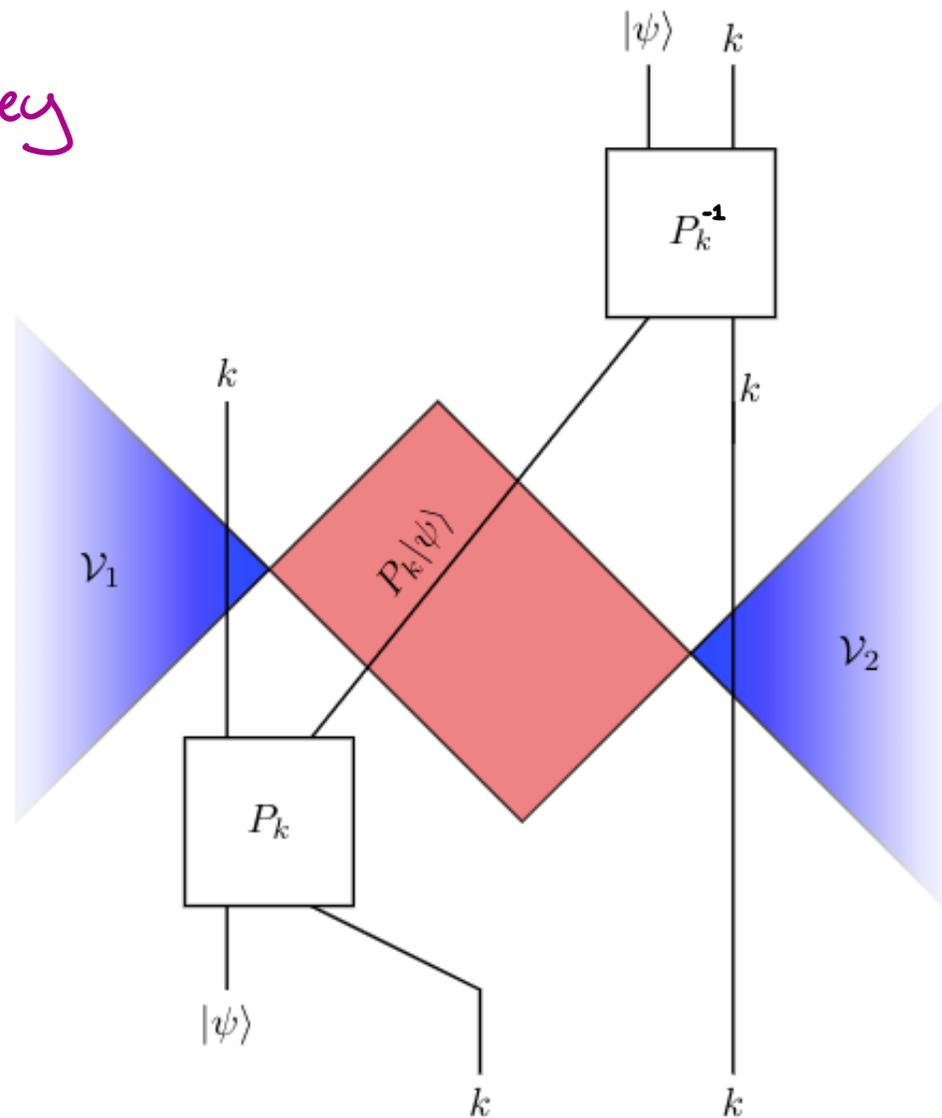
How to hide the state on Q

- Use the "one-time pad"

- 1) Prepare a string of bits, called k , the key and make a copy k'
- 2) Send k through V_2
- 3) Encode $|\psi\rangle$ using k' (do $P_{k'}$)
- 4) Send "encoded" Q through U
- 5) Undo P_k using k , recover $|\psi\rangle$

- Can show any such procedure must have $I(V_1:V_2|U)$ large

$$P_U \propto \sum_k P_k |\psi\rangle\langle\psi| P_k = \mathbb{I}$$



(Improved)

Connected wedge
theorem

Connected wedge theorem

- Pick four wedges, define:

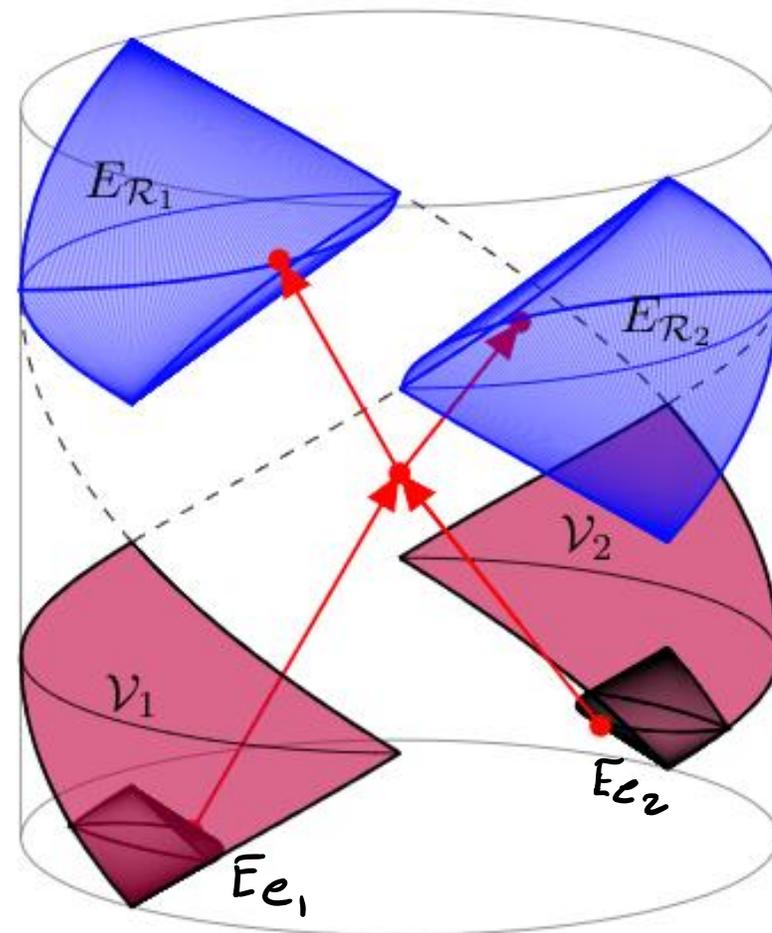
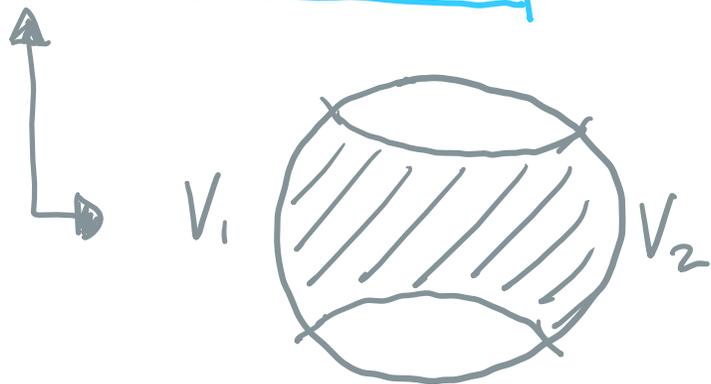
$$\mathcal{J}_{12 \rightarrow 12}^E \equiv \mathcal{J}^+(E_{e_1}) \cap \mathcal{J}^+(E_{e_2}) \cap \mathcal{J}^-(E_{r_1}) \cap \mathcal{J}^-(E_{r_2})$$

And:

$$V_1 \equiv \hat{\mathcal{J}}^+(e_1) \cap \hat{\mathcal{J}}^-(r_1) \cap \hat{\mathcal{J}}^-(r_2)$$

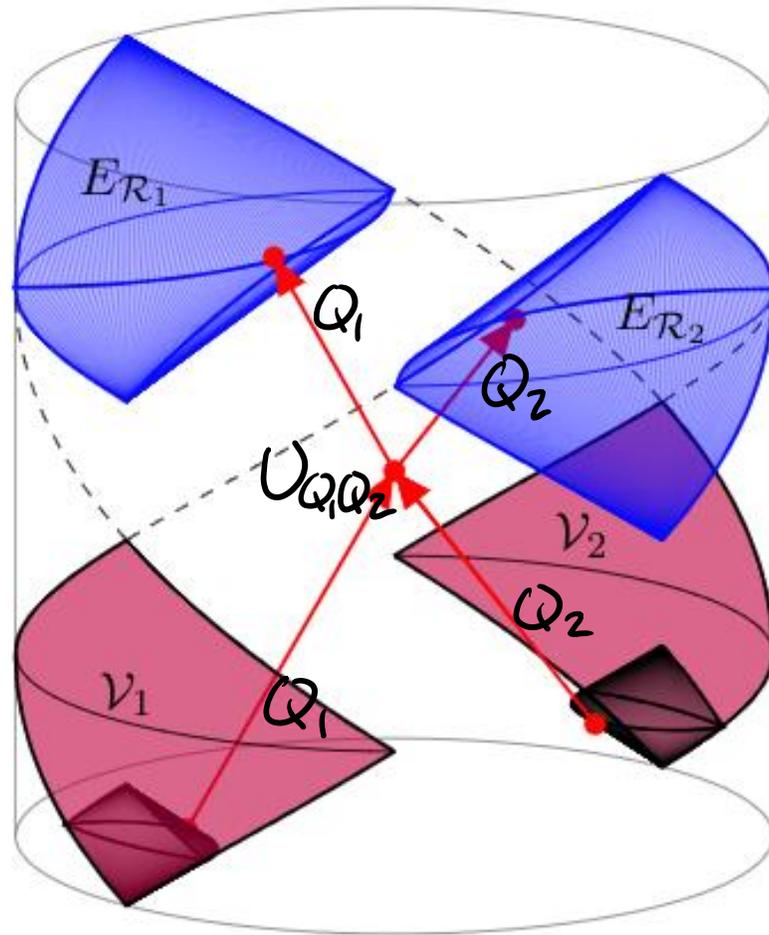
$$V_2 \equiv \hat{\mathcal{J}}^+(e_2) \cap \hat{\mathcal{J}}^-(r_1) \cap \hat{\mathcal{J}}^-(r_2)$$

If $\mathcal{J}_{12 \rightarrow 12}^E$ is non-empty,
then $I(V_1; V_2) = O(1/GN)$



Connected wedge theorem

Use same strategy...



For the boundary to reproduce the bulk unitary U_{Q_1, Q_2}
 V_1 and V_2 must share lots of entanglement $\rightarrow I(V_1; V_2) = \mathcal{O}(1/G_N)$

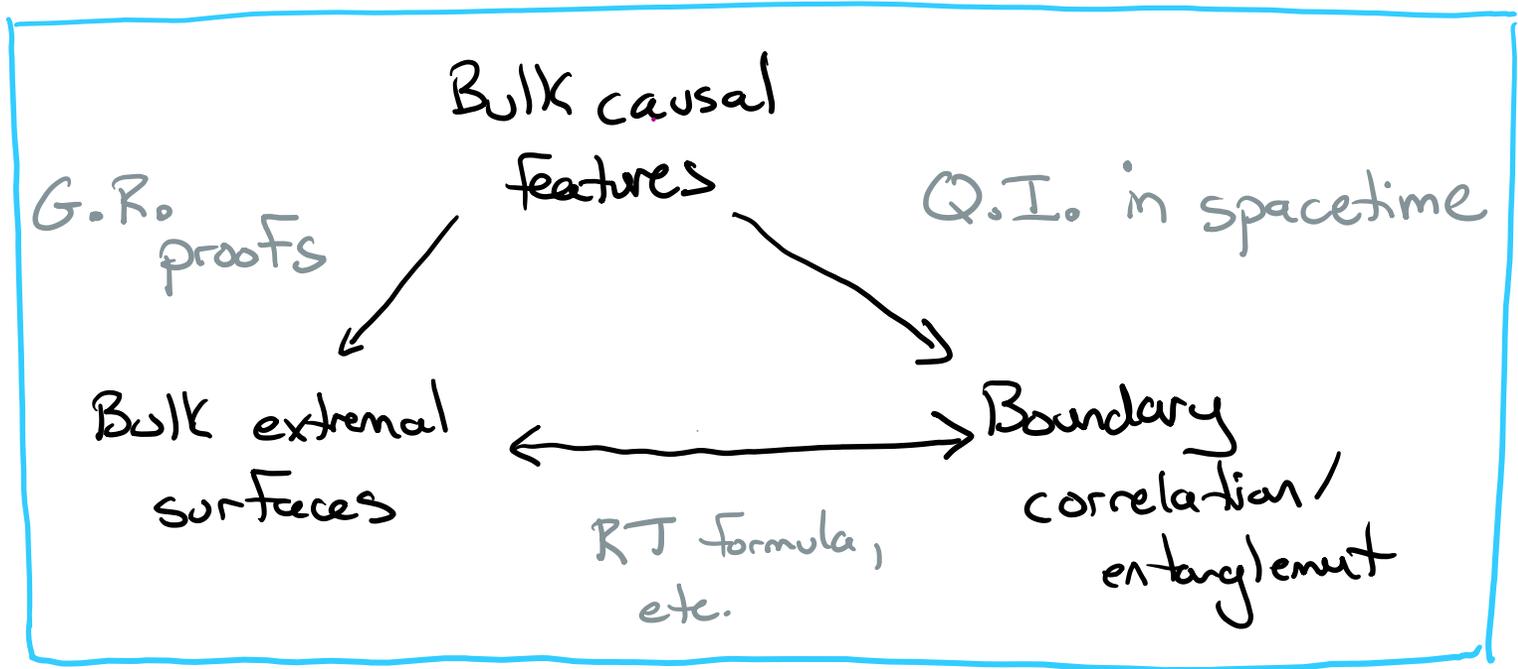
Final Remarks

Commut

- Two theorems together suggest a general causal feature \sim entanglement connection
- Complements usual extremal surface \sim entanglement connection given by RT

- Interesting to revisit problems studied using extremal surfaces, and look for a causal perspective:

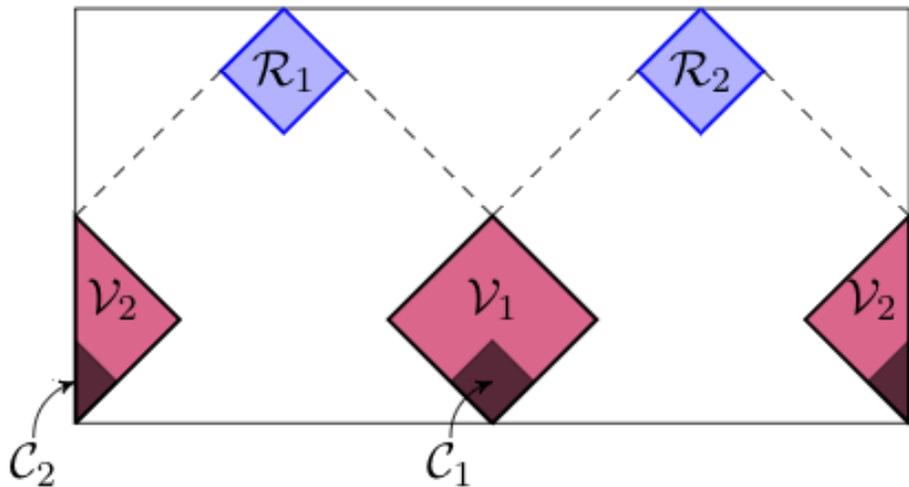
↳ e.g. In the BH + islands context, found causal condition for island formation (at least in toy models) (2102.01810 AM, David Wakeham)



Thanks! Questions?
(also, available on Slack
or may@phas.ubc.ca)

Improved connected wedge theorem

- Pick c_1, c_2, R_1, R_2 :



Define: $V_1 \equiv \hat{J}^+(c_1) \cap \hat{J}^-(R_1) \cap \hat{J}^-(R_2)$

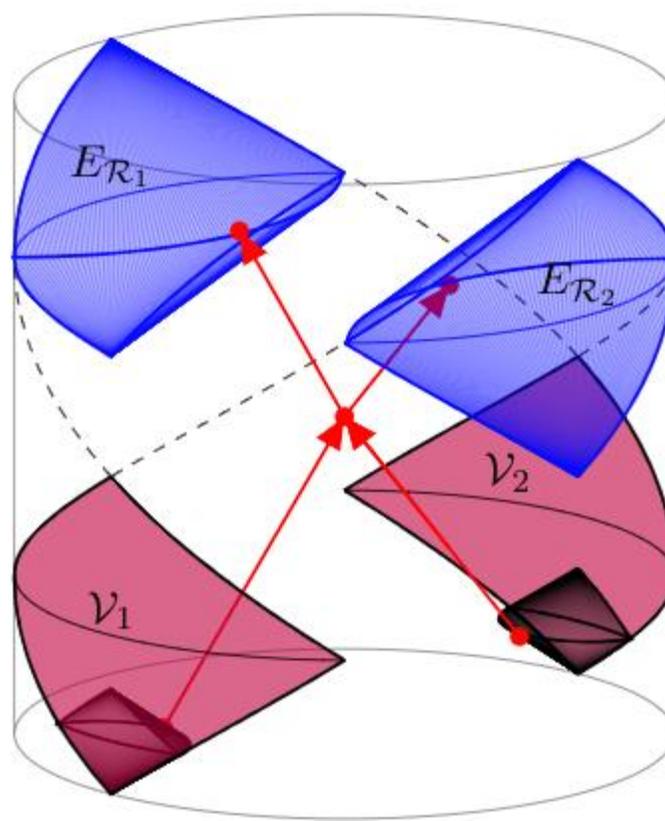
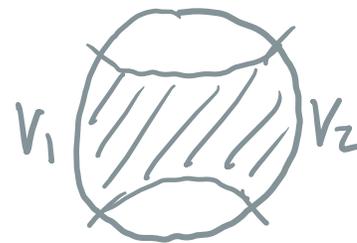
$V_2 \equiv \hat{J}^+(c_2) \cap \hat{J}^-(R_1) \cap \hat{J}^-(R_2)$,

$\mathcal{J}_{12 \rightarrow 12}^E \equiv \mathcal{J}^+(E_{c_1}) \cap \mathcal{J}^+(E_{c_2}) \cap \mathcal{J}^-(E_{R_1}) \cap \mathcal{J}^-(E_{R_2})$

Then:

$\mathcal{J}_{12 \rightarrow 12}^E \neq \emptyset$
 \Downarrow
 $\mathcal{I}(V_1:V_2) = O(1/G_N)$

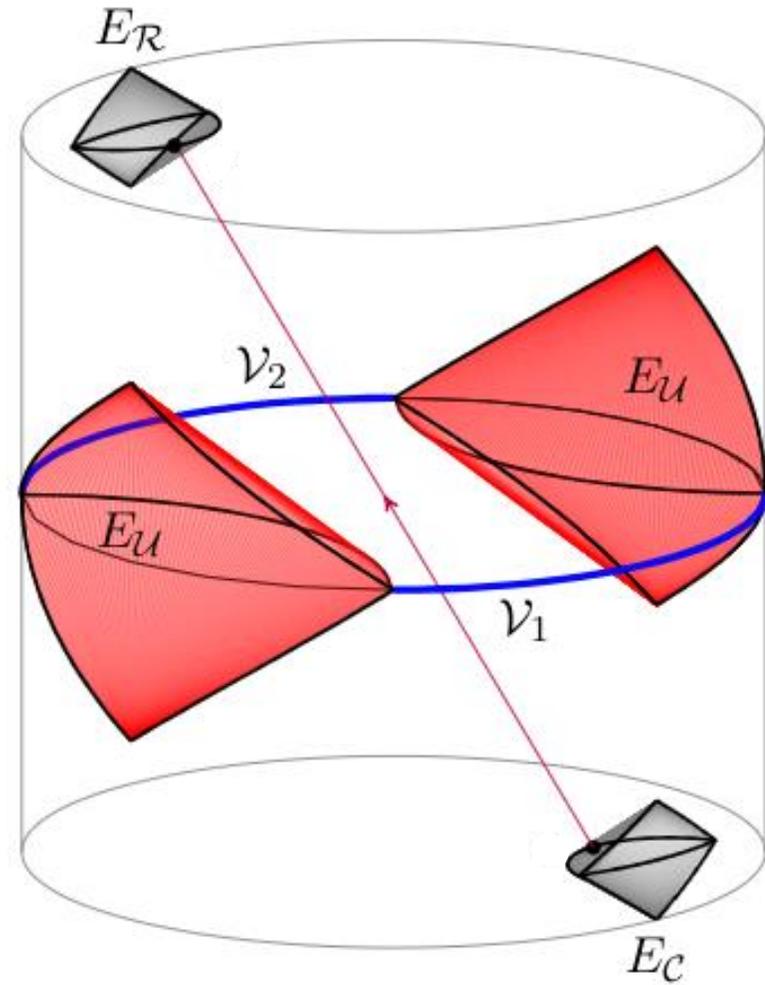
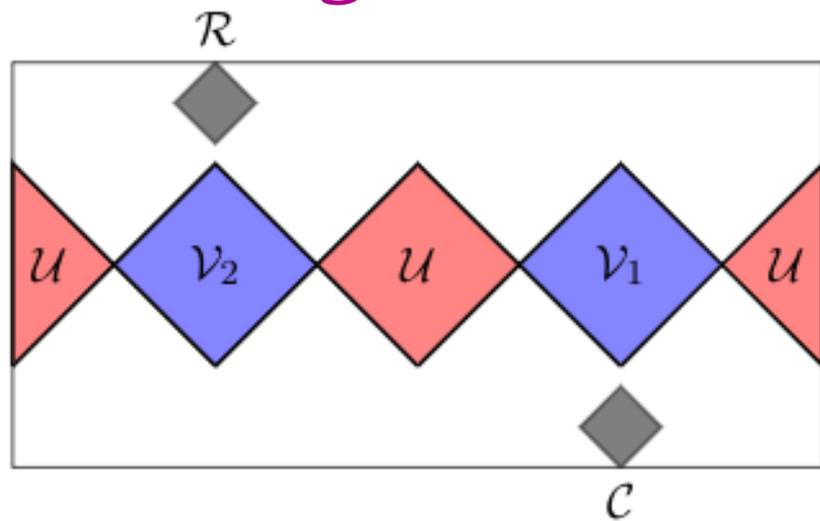
ie $E_{V_1 V_2} \neq E_{V_1} \cup E_{V_2}$,



"entanglement scattering region"

Privacy-Duality theorem

- Pick e, R, U in boundary, and a Cauchy surface Σ "through" U



Define: $V_1 \equiv \hat{D}(J^+(e) \cap \Sigma \cap U')$

$V_2 \equiv \hat{D}(J^-(R) \cap \Sigma \cap U')$

Then: IF \exists a causal curve from E_e to E_R that avoids E_U , then $I(V_1:V_2|U) = \alpha(1/\delta n)$