

# Moduli spaces, instantons, monopoles and Quantum Algebras

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# 4d gauge theories

- What can we do beyond perturbation theory?
- Are there hidden algebraic structures?
- What are the exactly computable quantities?

# 2d CFT

integrable lattice models  
quantum integrable systems

- conserved quantities
- holomorphic factorization
- conformal blocks, OPE
- Kac-Moody, Virasoro, W-algebra
- quantum groups, R-matrix

# What of these structures appear in gauge theories?

For 4d  $\mathcal{N} = 2$  theories

- vacua sector = integrable system
- holomorphic factorization
- emergence of quantum algebras
  - CFT type (Virasoro, affine Kac-Moody)
  - Spin chain type (R-matrix, quantum group)

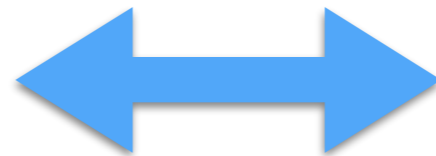
# What of these structures appear in gauge theories?

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# vacua = integrable system

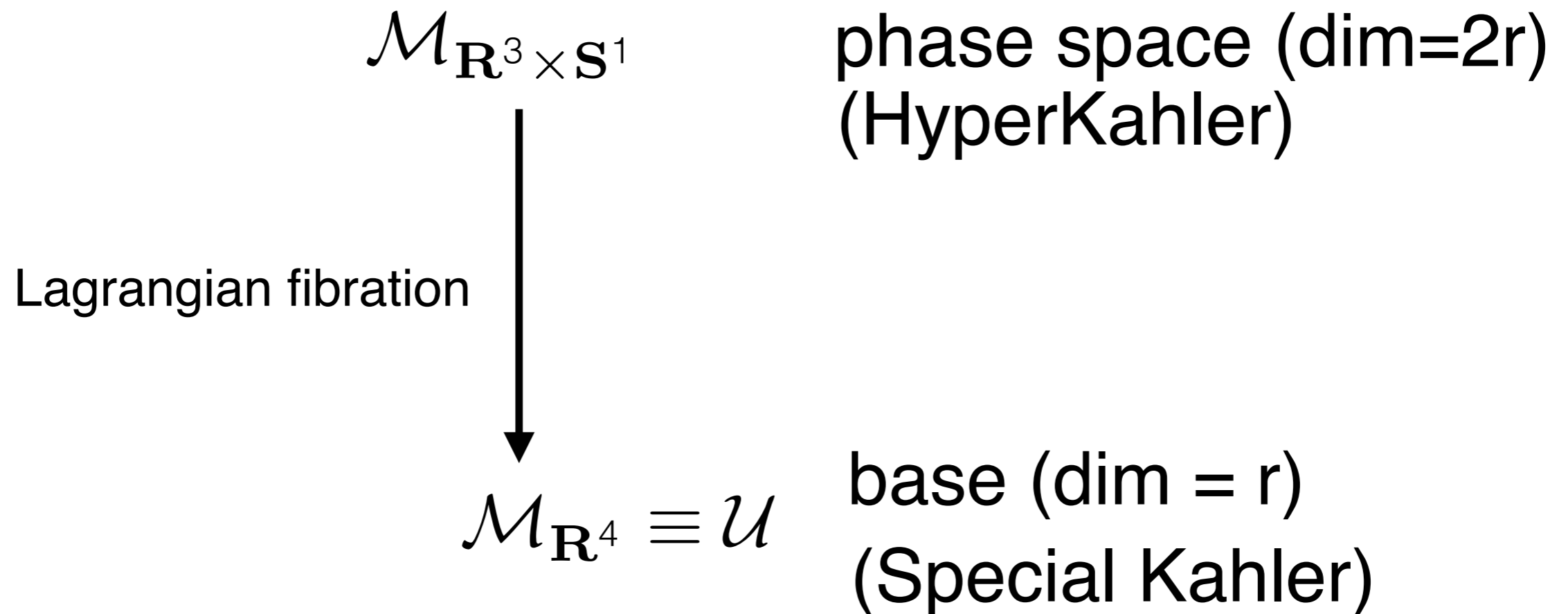
Coulomb branch of vacua  
of 4d  $\mathcal{N} = 2$  gauge theory



base of algebraic  
integrable system

**Seiberg-Witten'94**  
**GKMMM'95**  
**Donagi-Witten'95**  
**Seiberg-Witten'96**

# vacua = integrable system



# What we know?

For 4d  $\mathcal{N} = 2$  theories

- vacua sector = integrable system
- holomorphic factorization
- emergence of quantum algebras
  - Kac-Moody/Virasoro/W-algebra type
  - Quantum groups (R-matrix, spin chains)



# holomorphic factorization

$$Z_{\mathbb{S}^4} = \int da Z(a) \overline{Z(a)} \quad \text{VP'07}$$

$Z(a)$  = partition function on  $\mathbf{R}_{\varepsilon_1, \varepsilon_2}^4$  **Nekrasov'02**

$Z_{\varepsilon_1, \varepsilon_2}(a)$  = 2d CFT conformal block **Alday-Gaiotto-Tachikawa'08**

# holomorphic factorization

*700 pages review book  
coming out to arXiv next week*

Localization techniques in quantum field theories

**Vasily Pestun, Maxim Zabzine,  
Francesco Benini, Tudor Dimofte,  
Thomas T. Dumitrescu,  
Kazuo Hosomichi, Seok Kim,  
Kimyeong Lee, Bruno Le Floch,  
Marcos Marino, Joseph A. Minahan,  
David R. Morrison, Sara Pasquetti,  
Jian Qiu, Leonardo Rastelli,  
Shlomo S. Razamat, Silvu S. Pufu,  
Yuji Tachikawa, Brian Willett,  
Konstantin Zarembo**

$Z_{\varepsilon_1, \varepsilon_2}(a)$  is conformal block of what algebra?

Where does it come from?

# emergence of quantum algebras

If 4d theory of ‘class  $S_g(C)$ ’ then

$Z(a)$  = conformal block of  $W(\mathfrak{g})$  algebra  
on 2d surface  $C$

Alday-Gaiotto-Tachikawa’08

$$W(\mathfrak{sl}_2) = \text{Vir}$$

‘Class  $S_g(C)$ ’ is the 4d theory obtained by  
reduction of 6d (0,2) type  $\mathfrak{g}$  on  $C$

$$\mathcal{M}_{\mathbf{R}^3 \times S^1} = \mathcal{M}_{\text{Hitchin}}(C, \mathfrak{g})$$

# emergence of quantum algebras

$$\mathcal{M}_{\text{Hitchin}}(\mathcal{C}, \mathfrak{g}) \longrightarrow W_{\varepsilon_1, \varepsilon_2}(\mathfrak{g}) \text{ algebra}$$

**Alday-Gaiotto-Tachikawa '08**

**Nekrasov-Witten '12**

**Teschner '13**

# emergence of quantum algebras

$$\mathcal{M}_{\text{Hitchin}}(\mathcal{C}, \mathfrak{g}) \longrightarrow W_{\varepsilon_1, \varepsilon_2}(\mathfrak{g}) \text{ algebra}$$

**Alday-Gaiotto-Tachikawa '08**  
**Nekrasov-Witten '12**  
**Teschner '13**

What replaces  $W(\mathfrak{g})$ -algebra for generic  $\mathcal{M}$  ?

# emergence of quantum algebras

## *Proposition*

For generic  $\mathcal{M}$ , W-algebra  $W_{\varepsilon_1, \varepsilon_2}(\mathcal{M})$  is  $\varepsilon_1$ -quantized algebra of holomorphic functions on  $\mathcal{M}$  in HyperKähler  $\varepsilon_2/\varepsilon_1$ -rotated complex structure

$$W_{\varepsilon_1, \varepsilon_2}(\mathcal{M}) = \mathbb{C}_{\varepsilon_1}[\mathcal{M}_{\varepsilon_2/\varepsilon_1}]$$

Quantization  
(Planck constant)

HyperKähler  
twistor space  
parameter

**Nekrasov-Witten, Bellinson-Drinfeld**  
**Teschner, Kashaev**  
**Kontsevich-Soibelman, Gaiotto-Moore-Neitzke**  
**Fock-Goncharov,**  
**Neitzke, Vafa, Cecotti-Neitzke-Vafa**  
**Nekrasov-Rosly-Shatashvili, Gaiotto**  
**Kontsevich, Gukov-Witten**

# Quantization on HK

$(X, \omega)$  real symplectic manifold

Quantization has **one** parameter:

$\hbar$

Dirac  
Fedosov  
Kontsevich

$$\hat{f} * \hat{g} - \hat{g} * \hat{f} = \hbar \{f, g\}_{\omega^{-1}} + \dots$$

$(X, \omega_I, \omega_J, \omega_K)$  HK manifold

Quantization has **two** complex parameters:

1)  $\zeta \in \mathbb{C}\mathbb{P}^1$  twistor

2)  $\varepsilon \in \mathbb{C}$

Kontsevich  
Cattaneo-Felder  
Kapustin-Orlov  
Gukov-Witten

$(X, \hbar^{-1} \Omega_{\zeta})$  holomorphic symplectic

quantized by non-commutative open B-model



What we know about  $W_{\varepsilon_1, \varepsilon_2}(\mathcal{M})$  ?

associative  
algebra

## 1. Seiberg-Witten limit

$$\varepsilon_1 \rightarrow 0$$

$$\varepsilon_2 \rightarrow 0$$

$$W_{\varepsilon_1, \varepsilon_2}(\mathcal{M}) = \mathbb{C}[\mathcal{U}]$$

classical  
commutative  
ring

‘functions on u-plane’

‘Coulomb chiral ring in 4d’

‘classical commuting Hamiltonians’

## 2. Nekrasov-Shatashvili limit

$$\varepsilon_2 \rightarrow 0$$

$$\varepsilon_1 \rightarrow \hbar$$

$$W_{\varepsilon_1, \varepsilon_2}(\mathcal{M})$$

‘Ring of quantum commuting Hamiltonians’

quantum  
commutative  
ring

*The eigenvalues of elements in  $W_{\varepsilon_1, \varepsilon_2}(\mathcal{M})$  is spectrum of quantum integrable system*

What can we say  $W_{\varepsilon_1, \varepsilon_2}(\mathcal{M})$  for generic  $\varepsilon_1, \varepsilon_2$  ?

Two large classes of 4d theories

1. **class  $S_g$**  6d (0,2)  $g$ =ADE reduced on  $C$

$\mathcal{M}$  is  $g$ -type Hitchin system on  $C$



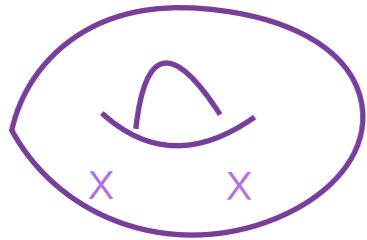
Witten'96  
Gaiotto'07  
AGT'08  
GMN'09

2.  **$\Gamma$ -quiver** 6d quiver theories reduced on  $C^\vee$   
 $\mathcal{M}$  is  $G_\Gamma$ -monopole system on  $C \times S^1$  (2d flat)

Douglas-Moore'96  
Cherkis-Kapustin'00  
Katz-Mayer-Vafa'97  
Nekrasov-VP-(Shatashvili) '12,'13  
Aganagic'15  
Nekrasov'15, Kimura-VP'15

# class $S_g$ theories

6d (0,2)  $\mathfrak{g}$ =ADE reduced on  $C$   
 $\mathcal{M}$  is  $\mathfrak{g}$ -Hitchin system on  $C$



$$\mathcal{M}_{\zeta \neq 0, \infty} \simeq \mathcal{M}_{\text{flat}}(C, G)$$



$W_{\varepsilon_1, \varepsilon_2}(\mathcal{M})$  is quantum cluster algebra  
(algebra of Verlinde loop operators)

Cheng's talk

Fock-Goncharov  
Teschner  
Gaiotto-Moore-Neitzke  
Alday-Gaiotto-Tachikawa

For punctured disk  $C = \mathbb{C}^\times$  and parabolic reduction

$W_{\varepsilon_1, \varepsilon_2}(\mathcal{M})$  is  $W(\mathfrak{g})$  algebra

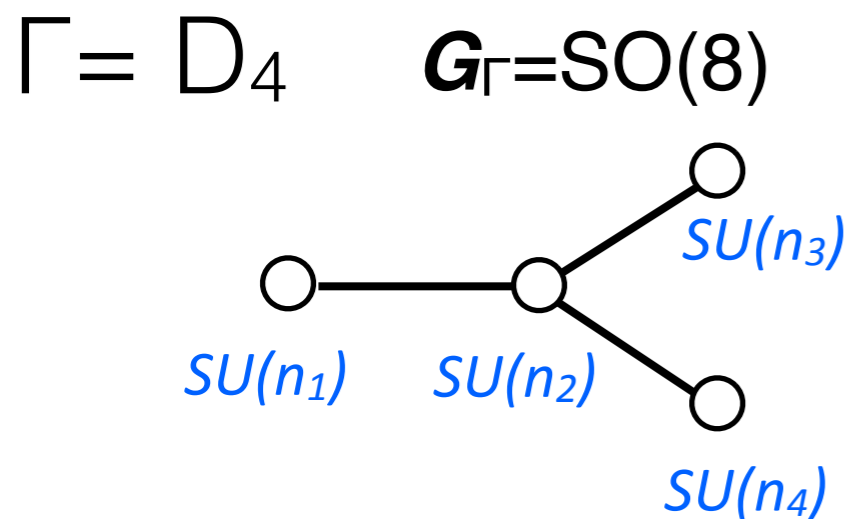
(integrals of motion in Toda 2d field theory)

Drinfeld-Sokolov  
Zamolodchikov  
Feigin-Frenkel  
Belinson-Drinfeld

# Quiver theories

6d  $\Gamma$ -quiver gauge theories reduced on  $\mathbb{C}^\vee$

$\mathcal{M}$  is  $G_\Gamma$ -monopole system on  $\mathbb{C} \times S^1$   
 (group valued analogue of Hitchin system)



$x \in \mathbb{C}$

$$\langle \chi_i[Y] \rangle = T_{n_i}(x) \text{-polynomial}$$

$\varepsilon_1 = 0, \varepsilon_2 = 0$   
 Nekrasov-VP '12

$\varepsilon_1 \neq 0, \varepsilon_2 = 0$   
 Nekrasov-VP-Shatashvili '13

$\varepsilon_1 \neq 0, \varepsilon_2 \neq 0$   
 Nekrasov '15, Kimura-VP '15

$$\chi_i[Y(x)] = \text{tr}_{L_i}(g(x))$$

$$[g(x)] = \prod_i y(x)^{a_i^\vee}$$

character

q-character =  $W_q$ -algebra ( $\mathfrak{g}_\Gamma$ )

Frenkel-Reshetikhin '98

qq-character =  $W_{qq}$ -algebra ( $\mathfrak{g}_\Gamma$ )

Frenkel-Reshetikhin '97

# UV-observables

$$Y_i(x) = \det(1 - x^{-1} e^{\Phi_i + \epsilon_1 J_1 + \epsilon_2 J_2})|_{\mathcal{E}_i}$$

qq-characters are proper IR-observables  
made from  $Y_i(x)$

**Nekrasov's talk**

# qq-characters

Nekrasov@String'14,16  
VP@String-Math'15

Example  $A_1$   $U(n)$   
 $\circ$

SU(2)-charge monopoles of charge  $n$  on  $C \times S^1$

good in UV

$$Y_i(x) = \det(1 - x^{-1} e^{\Phi_i + \epsilon_1 J_1 + \epsilon_2 J_2}) \Big|_{\mathcal{E}_i}$$



good in IR

$$\chi_1 = Y_1(x) + Y_1^{-1}(q^{-1}x)$$

$$q = q_1 q_2$$

$$q_i = e^{\epsilon_i}$$

1. classical limit ( $\varepsilon_1=\varepsilon_2=0$ )

$$\chi_1[y] = \text{tr} \begin{pmatrix} y(x) & 0 \\ 0 & y(x)^{-1} \end{pmatrix} = y(x) + \frac{1}{y(x)}$$

***q-character = T-polynomial***

**is classical SW spectral curve**

$$y + 1/y = T_n(x) = x^n + u_1 x^{n-1} + \cdots + u_n$$

**is spectral curve of SU(2) monopoles  
of charge  $n$  on  $\mathbb{C} \times S^1$**

2. quantum integrable system (NS-limit  $\varepsilon_1=\hbar$ ,  $\varepsilon_2=0$ )

$$Y(x) = Q(x)/Q(q^{-1}x)$$

*q-character = T-polynomial*

**is quantum SW spectral curve**

$$Q(qx) + Q(q^{-1}x) = Q(x)T(x)$$

(aka TQ-Baxter or analytic Bethe ansatz)

$$T_n(x) = x^n + u_1 x^{n-1} + \dots + u_n$$



**for generic quiver we find traces of R-matrices of rational, trigonometric and elliptic spin-chain types for any ADE and their affinization (toroidal)**

$$\text{Yang}(\mathfrak{g}), U_q(\hat{\mathfrak{g}}), E_{\mathfrak{g}, C}(\mathfrak{g})$$

Gerasimov-Kharchev-Lebedev-Oblezin'04

**The reason is that the phase space, the monopole moduli space  $M$  on  $C \times S^1$  is a symplectic leaf of the Poisson-Lie loop group  $\{\mathfrak{g}(x)\}$  with classical r-matrix Poisson bracket ! The quantization produces quantum R-matrix .**

$x \in C$  is spectral parameter

c.f. Costello's & Witten's talks

# q-characters=classical W-algebra

q-characters generate commutative

$W_q(\mathfrak{g})$  algebra

Nekrasov-VP-Shatashvili'13

Alternatively (after T-duality on the fibers of integrable system) this classical Poisson  $W_q(\mathfrak{g})$  algebra can be obtained by group version of Drinfeld-Sokolov reduction on group-valued  $q$ -'opers' ( $q$ -connections)

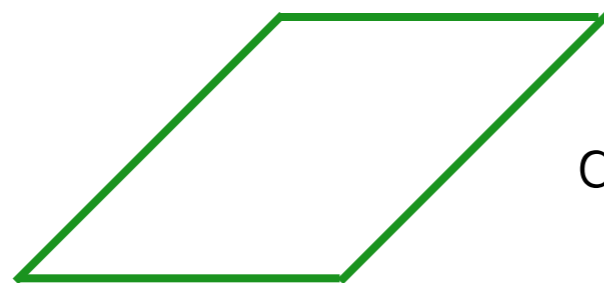
Steinberg'65

Semenov Tian Shansky'99

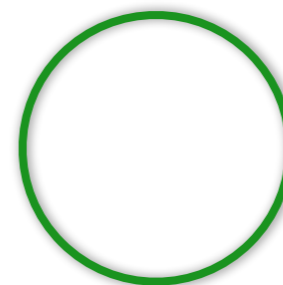
Sevostiyarov'00

$$\{g(x)\} / \{g(x) \sim s(x)g(x)s^{-1}(qx)\}$$

$C \times_q S^1$



q-twisted over



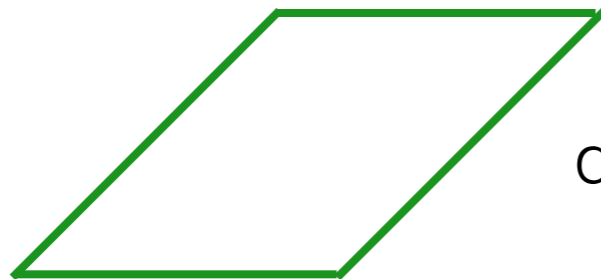
# qq-characters=quantum W-algebra

qq-characters generate associative  $W_{qq}(\mathfrak{g})$  algebra: quantization of  $W_q(\mathfrak{g})$

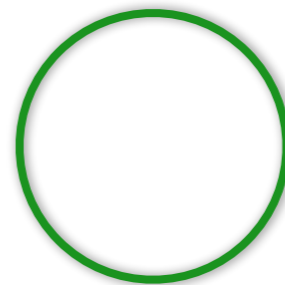
$$W_{qq}(\mathfrak{g}) = \mathbb{C}_{\epsilon_1} [\mathcal{M}_{\text{mono}}^{G,n}(C \times_{\epsilon_2} S^1)]$$

$$\{g(x)\} / \{g(x) \sim s(x)g(x)s^{-1}(qx)\}$$

$C \times_q S^1$



q-twisted over



# Free field realization of quiver $W_{\mathfrak{q}\mathfrak{q}}$ -algebra

Kimura-VP'15

Consider quiver gauge theory  
extended Nekrasov partition function

Nekrasov-  
Marshakov

$$Z(t) = \int_{\mathfrak{M}_\Gamma} \exp \left( \sum_i \sum_{p=1}^{\infty} t_{i,p} \mathbf{ch} Y_i^{[p]} \right)$$

which depends on infinite set of parameters

$$\{t\} = \{t_{i,p}\} \quad \begin{array}{l} p = 1, 2, \dots, \infty \\ i \in \Gamma_0 \end{array}$$

(generating function of intersection of all Chern  
characters of  $Y$ -bundle over instanton moduli space)

# Free field realisation of quiver $W_{\mathbf{qq}}$ -algebra

Interpret  $Z\{t\}$  is a state in Fock space  $\mathbb{C}[[t]]$   
of Heisenberg algebra generated by  $\left\{t_{i,p}, \frac{\partial}{\partial t_{i,p}}\right\}$

**Kimura-VP'15:**  
c.f. Aganagic talk

$$|Z\rangle = \prod_{x \in \mathcal{X}} \hat{S}(x) |1\rangle$$

$$\hat{S}(x) = \int d_{q_2} x S(x)$$

$$S_{i,x} =: \exp \left( \sum_{p>0} s_{i,-p} x^p + s_{i,0} \log x + \tilde{s}_{i,0} + \sum_{p>0} s_{i,p} x^{-p} \right) :$$

Here the screening current

$$S_{i,x} =: \exp \left( \sum_{p>0} s_{i,-p} x^p + s_{i,0} \log x + \tilde{s}_{i,0} + \sum_{p>0} s_{i,p} x^{-p} \right) :$$

is realised in terms of the free fields:

$$s_{i,-p} \stackrel{p>0}{=} (1 - q_1^p) t_{i,p}, \quad s_{i,0} = t_{i,0}, \quad s_{i,p} \stackrel{p>0}{=} -\frac{1}{p} \frac{1}{1 - q_2^{-p}} c_{ij}^{[p]} \partial_{i,p}$$

$c_{ij}$  is deformed Cartan matrix of quiver

$$c_{ij} = \left( \delta_{ji} - \sum_{e:j \rightarrow i} \mu_e^{-1} \right) + \left( q^{-1} \delta_{ji} - \sum_{e:i \rightarrow j} q^{-1} \mu_e \right)$$

Elliptic version and non-simply laced version  
are defined similarly for generic<sup>30</sup> quiver

**A<sub>1</sub> : Kozcaz's talk**

# qq-characters = free field realisation of $W_{q_1, q_2}(\mathfrak{g})$

*Example*  $A_1 \overset{\text{U}(n)}{\circ}$

$$T_1(x) = Y_1(x) + Y_1^{-1}(q^{-1}x)$$

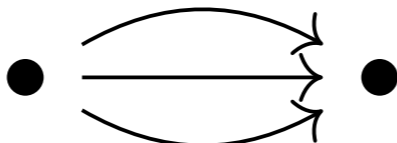
where

$$Y_{i,x} =: \exp \left( \sum_{p>0} y_{i,-p} x^p + y_{i,0} + \sum_{p>0} y_{i,p} x^{-p} \right) :$$

$$y_{i,-p} = (1 - q_1^p)(1 - q_2^p)(\tilde{c}^{[-p]})_{ij} t_{j,p}, \quad y_{i,0} = -\tilde{c}_{ij}^{[0]} t_{i,0} \log q_2 \quad y_{i,p} = -\frac{1}{p} \partial_{i,p}$$

$$\langle 1 | T_1(x) | Z \rangle = \text{polynomial deg } n$$

The construction of  $\mathbf{W}_{q_1, q_2}(\Gamma)$  in **Kimura-VP'15** is more general than  $\mathbf{W}_{q_1, q_2}(\mathfrak{g})$  in **Frenkel-Reshetikhin'97** and applicable to any quiver, not necessarily of finite Dynkin type. It depends on parameters  $\mu \in H^1(\Gamma, \mathbb{C}^\times)$

Hyperbolic example   $c = \begin{pmatrix} 2 & -3 \\ -3 & 2 \end{pmatrix}$

$$T_{1,x} = Y_{1,x} + q_1^{-1} : \frac{Y_{2,\mu_1^{-1}x} Y_{2,\mu_2^{-1}x} Y_{2,\mu_3^{-1}x}}{Y_{1,q^{-1}x}} : \quad \mathcal{S}(w) = \frac{(1 - q_1 w)(1 - q_2 w)}{(1 - qw)(1 - w)}$$

$$+ q_1^{-2} \left[ \mathcal{S}(\mu_1 \mu_2^{-2}) \mathcal{S}(\mu_1 \mu_3^{-1}) : \frac{Y_{1,\mu_1^{-1} \mu_2 q^{-1}} Y_{1,\mu_1^{-1} \mu_3 q^{-1}} Y_{2,\mu_2^{-1}x} Y_{2,\mu_3^{-1}x}}{Y_{2,\mu_1^{-1} q^{-1}x}} : \right.$$

$$+ \mathcal{S}(\mu_2 \mu_1^{-2}) \mathcal{S}(\mu_2 \mu_3^{-1}) : \frac{Y_{1,\mu_2^{-1} \mu_1 q^{-1}} Y_{1,\mu_2^{-1} \mu_3 q^{-1}} Y_{2,\mu_1^{-1}x} Y_{2,\mu_3^{-1}x}}{Y_{2,\mu_2^{-1} q^{-1}x}} : \left. \right.$$

$$+ \mathcal{S}(\mu_3 \mu_1^{-2}) \mathcal{S}(\mu_3 \mu_2^{-1}) : \frac{Y_{1,\mu_3^{-1} \mu_1 q^{-1}} Y_{1,\mu_3^{-1} \mu_2 q^{-1}} Y_{2,\mu_1^{-1}x} Y_{2,\mu_2^{-1}x}}{Y_{2,\mu_3^{-1} q^{-1}x}} : \left. \right]$$

$$+ \dots$$



$$\mathbf{Z}_{\text{gauge}} = \mathbf{Z}_{\text{top}}$$

W-algebra of ('refined' topological strings)

What are higher times  $\{t\}$  ?

Elliptic version and non-simply laced version  
are defined similarly for generic quiver

**A<sub>1</sub> Kozcaz's talk**

Generalization of Nakajima's results on equivariant K-theory of quiver variety and representation theory of  $U_q(\mathfrak{Lg})$  to affine Kac-Moody  $\mathfrak{g}$  with non-zero central charge

Quantum  $q$ -geometric Langlands is transparent and completely geometric:  
exchange of  $\varepsilon_1$  vs  $\varepsilon_2$

Rich algebraic structures coming from quantization of HyperKähler moduli spaces

Thanks to my collaborators

**N. Nekrasov**  
**S. Shatashvili**  
**T. Kimura**

*Thank you*