# Applications of AdS/CFT to elementary particle and condensed matter physics

# Johanna Erdmenger

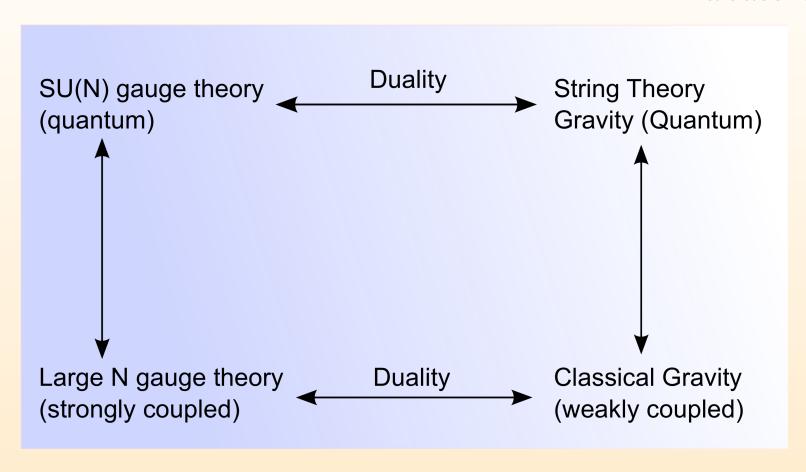
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### Starting point: AdS/CFT correspondence

#### Maldacena 1997



$$N \to \infty \Leftrightarrow g_s \to 0$$

't Hooft coupling  $\lambda$  large  $\Leftrightarrow \alpha' \to 0$ , energies kept fixed

Conjecture extends to more general gravity solutions

 $AdS_n \times S^m$  generalizes to more involved geometries

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Dual also to non-conformal, non-supersymmetric field theories

Gauge/gravity duality

- Important approach to studying strongly coupled systems
- New links of string theory to other areas of physics

# Gauge/gravity duality

#### QCD:

- Quark-gluon plasma
- Lattice gauge theory
- External magnetic fields

#### Condensed matter:

- Quantum phase transitions
- Conductivities and transport processes
- Holographic superconductors
- Kondo model, Weyl semimetals

Universality

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Renormalization group:

Large-scale behaviour is independent of microscopic degrees of freedom

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Large-scale behaviour is independent of microscopic degrees of freedom

 The same physical phenomenon may occur in different branches of physics

# Introduction: Top-down and bottom-up approach

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- a) Ten- or eleven-dimensional (super-)gravity
- b) Probe branes

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- Many examples where dual field theory Lagrangian is known

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- b) Probe branes
- Few parameters
- Many examples where dual field theory Lagrangian is known

#### Bottom-up approach:

Choose simpler, mostly four- or five-dimensional gravity actions

QCD: Karch, Katz, Son, Stephanov; Pomerol, Da Rold; Brodsky, De Teramond; ........

Condensed matter: Hartnoll et al, Herzog et al, Schalm, Zaanen et al, McGreevy, Liu, Faulkner et al; Sachdev et al .......

### Outline

- 1. Kondo effect
- 2. Condensation to new ground states; external magnetic field
- 3. Mesons
- 4. Axial anomaly

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Unifying theme: Universality



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1. Kondo model: Simple model for a RG flow with dynamical scale generation

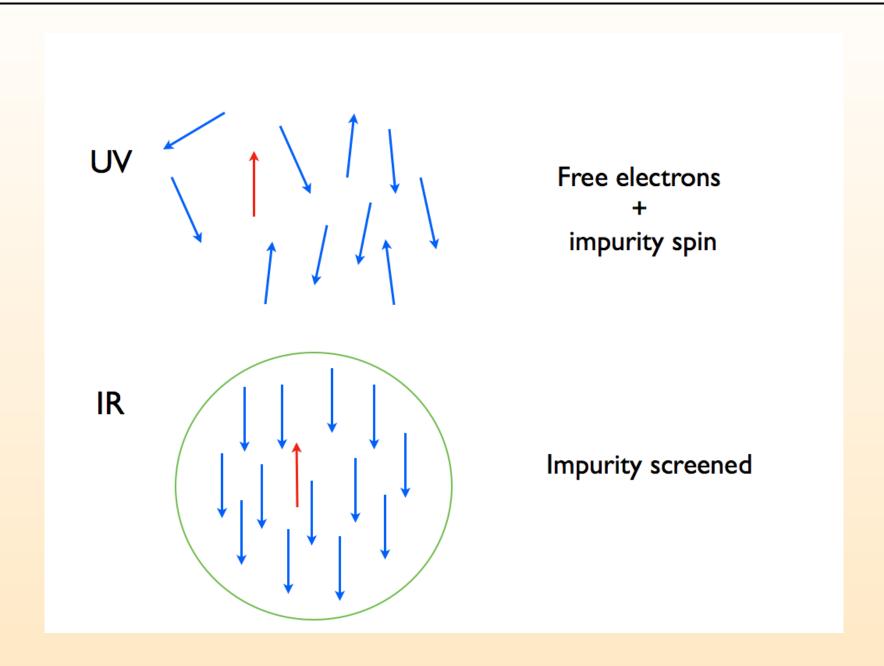
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Screening of a magnetic impurity by conduction electrons at low temperatures

Motivation for study within gauge/gravity duality:

- 1. Kondo model: Simple model for a RG flow with dynamical scale generation
- 2. New applications of gauge/gravity duality to condensed matter physics

# Kondo effect



Original Kondo model (Kondo 1964): Magnetic impurity interacting with free electron gas

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Logarithmic rise of conductivity at low temperatures

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Due to symmetries: Model effectively (1+1)-dimensional

Hamiltonian:

$$H = \frac{v_F}{2\pi} \psi^{\dagger} i \partial_x \psi + \lambda_K v_F \delta(x) \vec{S} \cdot \vec{J}, \quad \vec{J} = \psi^{\dagger} \frac{1}{2} \vec{T} \psi$$

Decisive in development of renormalization group IR fixed point, CFT approach Affleck, Ludwig '90's

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In this case, interaction term simplifies introducing slave fermions:

$$S^a = \chi^{\dagger} T^a \chi$$

Totally antisymmetric representation: Young tableau with Q boxes

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Screened phase has condensate  $\langle \mathcal{O} \rangle$ 

Parcollet, Georges, Kotliar, Sengupta cond-mat/9711192 Senthil, Sachdev, Vojta cond-mat/0209144

### Previous studies of holographic models with impurities:

Supersymmetric defects with localized fermions

Kachru, Karch, Yaida; Harrison, Kachru, Torroba Jensen, Kachru, Karch, Polchinski, Silverstein

Benincasa, Ramallo; Itsios, Sfetsos, Zoakos; Karaiskos, Sfetsos, Tsatis

Mück; Faraggi, Pando Zayas; Faraggi, Mück, Pando Zayas

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### Here: Model describing an RG flow

J.E., Hoyos, O'Bannon, Wu 1310.3271, JHEP 1312 (2013) 086

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Coupling of a magnetic impurity to a strongly interacting non-Fermi liquid

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#### Results:

- RG flow from perturbation by 'double-trace' operator
- Dynamical scale generation
- $AdS_2$  holographic superconductor
- Power-law scaling of conductivity in IR with real exponent
- Screening, phase shift

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#### Results:

- RG flow from perturbation by 'double-trace' operator
- Dynamical scale generation
- AdS<sub>2</sub> holographic superconductor
- Power-law scaling of conductivity in IR with real exponent
- Screening, phase shift
- Generalizations: Quantum quenches, Kondo lattices

J.E., Hoyos, O'Bannon, Wu 1310.3271, JHEP 1312 (2013) 086

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### Top-down brane realization

	0	1	2	3	4	5	6	7	8	9
N D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X
$N_5$ D5	X				X	X	X	X	X	

- 3-7 strings: Chiral fermions  $\psi$  in 1+1 dimensions
- 3-5 strings: Slave fermions  $\chi$  in 0+1 dimensions
- 5-7 strings: Scalar (tachyon)

# Near-horizon limit and field-operator map

D3:  $AdS_5 \times S^5$ 

D7:  $AdS_3 \times S^5 \to \text{Chern-Simons } A_{\mu} \text{ dual to } J^{\mu} = \psi^{\dagger} \sigma^{\mu} \psi$ 

D5:  $AdS_2 \times S^4 \rightarrow \left\{ \begin{array}{l} \mathsf{YM} \, a_t \, \mathsf{dual} \, \mathsf{to} \, \chi^\dagger \chi = q \\ \mathsf{Scalar} \, \mathsf{dual} \, \mathsf{to} \, \psi^\dagger \chi \end{array} \right.$ 

Operator		Gravity field
Electron current $J$	$\Leftrightarrow$	Chern-Simons gauge field $A$ in $AdS_3$
Charge $q = \chi^{\dagger} \chi$	$\Leftrightarrow$	2d gauge field $a$ in $AdS_2$
Operator $\mathcal{O} = \psi^{\dagger} \chi$	$\Leftrightarrow$	2d complex scalar $\Phi$

### Bottom-up model

#### Action:

$$S = S_{CS} + S_{AdS_2}$$

$$S_{CS} = -\frac{N}{4\pi} \int \text{Tr} \left( A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right),$$

$$S_{AdS_2} = -N \int d^3x \, \delta(x) \sqrt{-g} \left[ \frac{1}{4} \text{Tr} f^{mn} f_{mn} + g^{mn} \left( D_m \Phi \right)^{\dagger} D_n \Phi + V(\Phi^{\dagger} \Phi) \right],$$

$$D_{\mu} \Phi = \partial_m \Phi + i A_{\mu} \Phi - i a_{\mu} \Phi$$

#### Metric: BTZ black hole

$$ds^{2} = g_{\mu\nu}dx^{\mu}dx^{\nu} = \frac{1}{z^{2}} \left( \frac{dz^{2}}{h(z)} - h(z) dt^{2} + dx^{2} \right), \qquad h(z) = 1 - z^{2}/z_{H}^{2}$$
$$T = 1/(2\pi z_{H})$$

# 'Double-trace' deformation by $\mathcal{OO}^{\dagger}$

### Boundary expansion

$$\Phi = z^{1/2}(\alpha \ln z + \beta)$$
$$\alpha = \kappa \beta$$

 $\kappa$  dual to double-trace deformation

Witten hep-th/0112258

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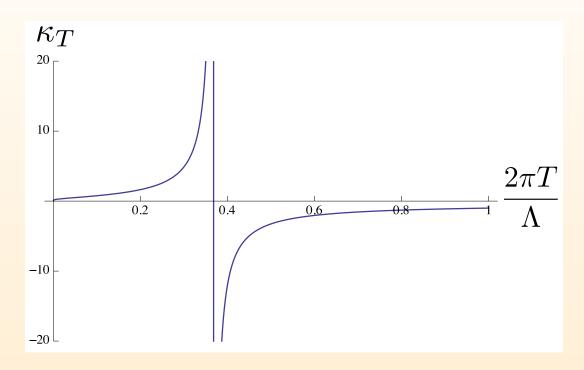
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### Dynamical scale generation

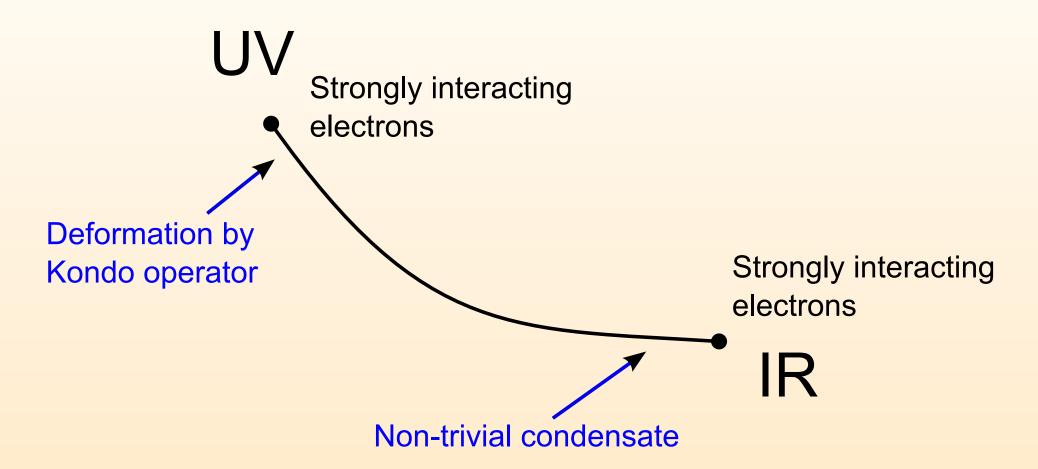
### Scale generation



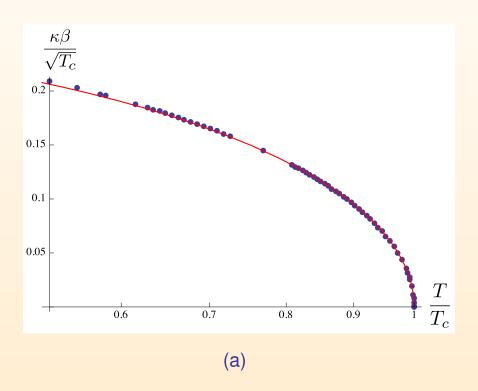
Divergence of Kondo coupling determines Kondo temperature

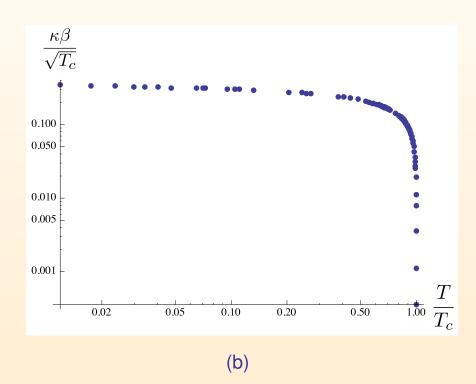
Below this temperature, scalar condenses





Normalized condensate  $\langle \mathcal{O} \rangle \equiv \kappa \beta$  as function of the temperature

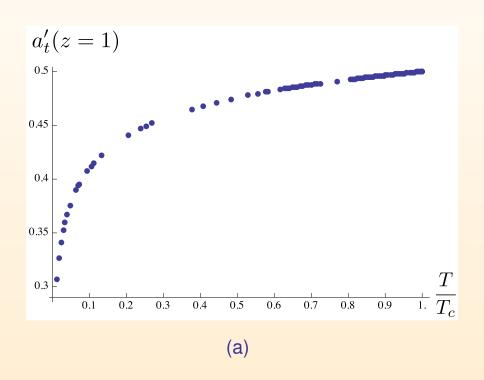


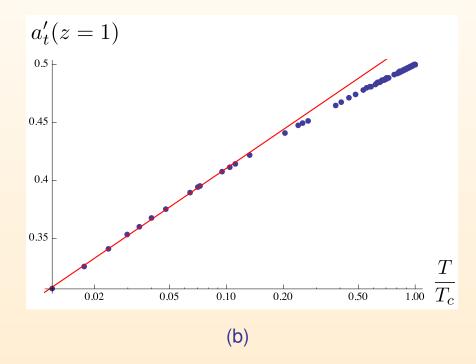


Mean field transition

 $\langle \mathcal{O} \rangle$  approaches constant for  $T \to 0$ 

#### Electric flux at horizon

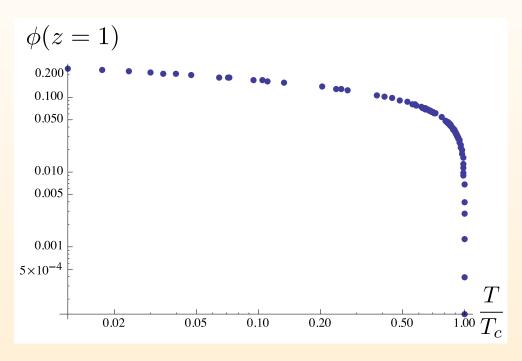




$$\sqrt{-g}f^{tr}\Big|_{\partial AdS_2} = q = \chi^{\dagger}\chi$$

# Impurity is screened

Resistivity from leading irrelevant operator (No log behaviour due to strong coupling)



#### IR fixed point stable:

Flow near fixed point governed by operator dual to 2d YM-field  $a_t$ 

$$\Delta = \frac{1}{2} + \sqrt{\frac{1}{4} + 2\phi_{\infty}^2}, \qquad \phi(z=1) = \phi_{\infty}$$

# Resistivity from leading irrelevant operator

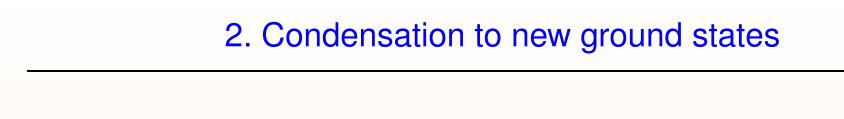
Entropy density:  $s=s_0+c_s\lambda_{\mathcal{O}}^2T^{-2+2\Delta}$ Resistivity:  $\rho=\rho_0+c_+\lambda_{\mathcal{O}}^2T^{-1+2\Delta}$ 

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#### **Outlook:**

- Transport properties, thermodynamics; entanglement entropy
- Quench
- Kondo lattice



# 2. Condensation to new ground states

Starting point: Holographic superconductors

Gubser 0801.2977; Hartnoll, Herzog, Horowitz 0803.3295

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Gubser, Pufu 0805.2960; Roberts, Hartnoll 0805.3898

Triplet pairing

Condensate breaks rotational symmetry

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Probe brane model reveals that field-theory dual operator is similar to  $\rho$ -meson: Ammon, J.E., Kaminski, Kerner 0810.2316

$$\langle \bar{\psi}_u \gamma_\mu \psi_d + \bar{\psi}_d \gamma_\mu \psi_u + bosons \rangle$$

# p-wave holographic superconductor

Einstein-Yang-Mills-Theory with SU(2) gauge group

$$S = \int d^5 x \sqrt{-g} \left[ \frac{1}{2\kappa^2} (R - 2\Lambda) - \frac{1}{4\hat{g}^2} F^a_{\mu\nu} F^{a\mu\nu} \right]$$

$$\alpha = \frac{\kappa_5}{\hat{g}}$$

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#### Gauge field ansatz

$$A = \phi(r)\tau^{3}dt + w(r)\tau^{1}dx$$

$$\phi(r) \sim \mu + \dots \qquad w(r) \sim d/r^{2}$$

 $\mu$  isospin chemical potential, explicit breaking  $SU(2) \to U(1)_3$  condensate  $d \propto \langle J_x^1 \rangle$ , spontaneous symmetry breaking



### Universality: Shear viscosity over entropy density

# Transport properties

Universal result of AdS/CFT:

Kovtun, Policastro, Son, Starinets

$$\frac{\eta}{s} = \frac{1}{4\pi} \frac{\hbar}{k_B}$$

Shear viscosity/Entropy density

Proof of universality relies on isotropy of spacetime

Metric fluctuations ⇔ helicity two states

Rotational symmetry broken  $\Rightarrow$  shear viscosity becomes tensor

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### p-wave superconductor:

Fluctuations characterized by transformation properties under unbroken SO(2):

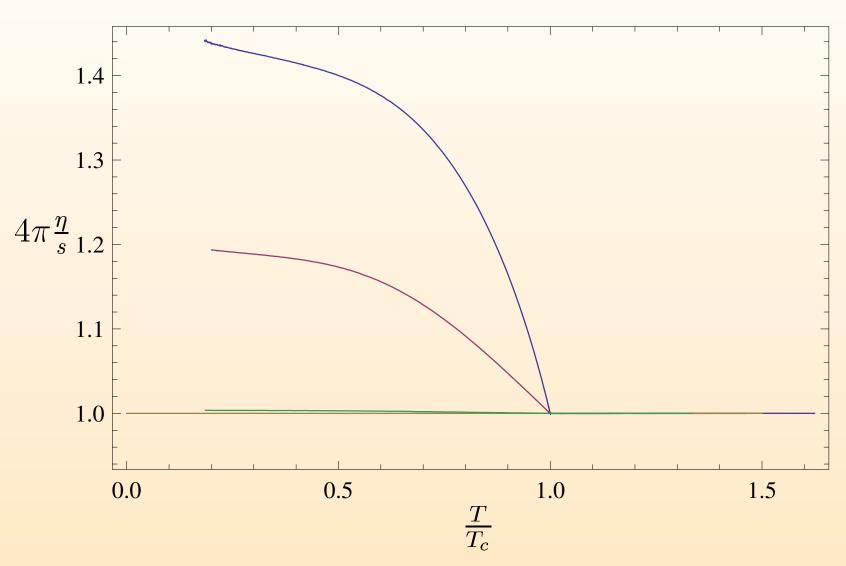
Condensate in *x*-direction:

 $h_{yz}$  helicity two,  $h_{xy}$  helicity one

J.E., Kerner, Zeller 1011.5912; 1110.0007

Backreaction: Ammon, J.E., Graß, Kerner, O'Bannon 0912.3515

J.E., Kerner, Zeller 1011.5912



 $\eta_{yz}/s=1/4\pi$ ;  $\eta_{xy}/s$  dependent on T and on  $\alpha$ 

Non-universal behaviour at leading order in  $\lambda$  and N

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Viscosity bound preserved  $\leftrightarrow$  Energy-momentum tensor remains spatially isotropic,  $T^{xx}=T^{yy}=T^{zz}$ 

Donos, Gauntlett 1306.4937

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Rebhan, Steineder 1110.6825

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Further recent anisotropic holographic superfluids:

Jain, Kundu, Sen, Sinha, Trivedi 1406.4874; Critelli, Finazzo, Zaniboni, Noronha 1406.6019

Recall: Necessary isospin chemical potential provided by non-trivial  $A_t^3(r)$ 

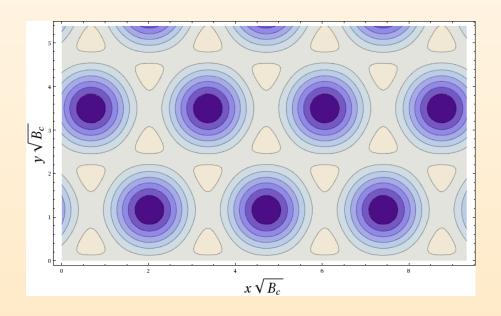
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#### For $B > B_c$ , the new ground state is a triangular lattice



Bu, J.E., Strydom, Shock 1210.6669

# External electromagnetic fields

A magnetic field leads to

 $\rho$  meson condensation and superconductivity in the QCD vacuum

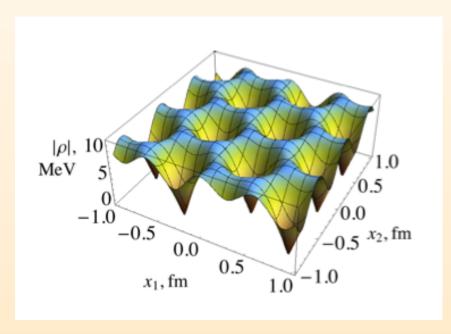
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#### Effective field theory:

Chernodub 1101.0117

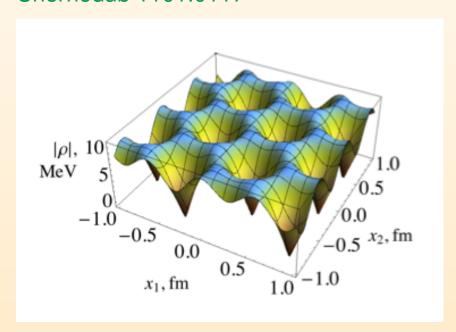


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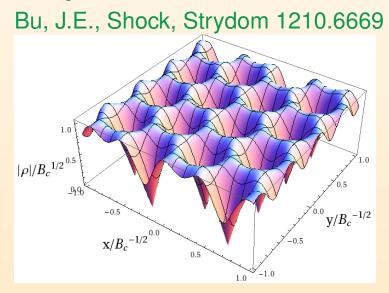
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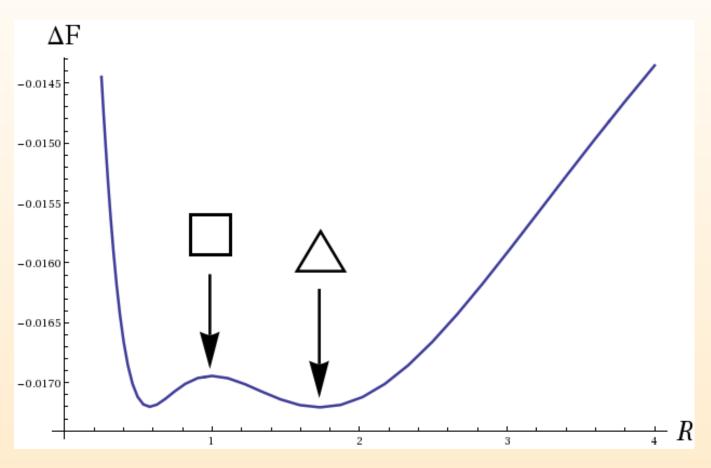


Gauge/gravity duality magnetic field in black hole supergravity background



Free energy as function of  $R = \frac{L_x}{L_y}$ 

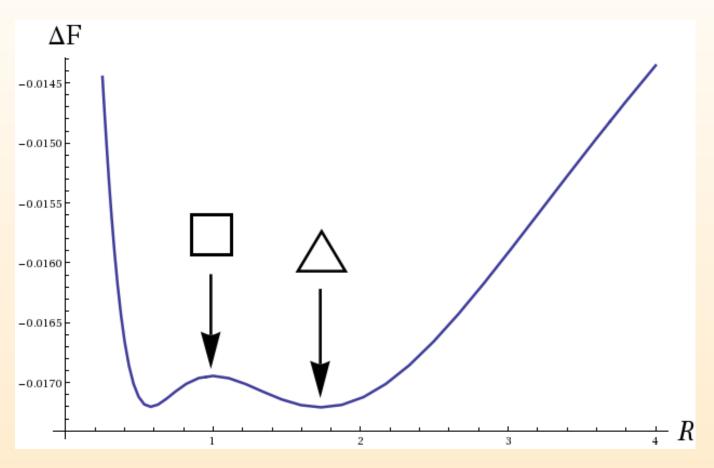
Bu, J.E., Shock, Strydom 1210.6669



# Free energy

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Bu, J.E., Shock, Strydom 1210.6669



Lattice generated dynamically

# Spontaneously generated lattice ground state in magnetic field

- Ambjorn, Nielsen, Olesen '80s: Gluon or W-boson instability
   Fermions: Z<sub>2</sub> topological insulator Beri, Tong, Wong 1305.2414
- Chernodub '11-'13: ho meson condensate in effective field theory, lattice Note:  $B_{
  m crit}\sim m_{
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Similar condensation in Sakai-Sugimoto model Callebaut, Dudas, Verschelde 1105.2217



# Spontaneously generated inhomogeneous ground states

### With magnetic field:

Bolognesi, Tong; Donos, Gauntlett, Pantelidou; Jokela, Lifschytz, Lippert; Cremonini, Sinkovics; Almuhairi, Polchinski.

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#### With Chern-Simons term at finite momentum:

Domokos, Harvey;

Helical phases: Nakamura, Ooguri, Park; Donos, Gauntlett

Charge density waves: Donos, Gauntlett; Withers;

Rozali, Smyth, Sorkin, Stang.

## 3. Quarks in the AdS/CFT correspondence

#### D7-Brane probes

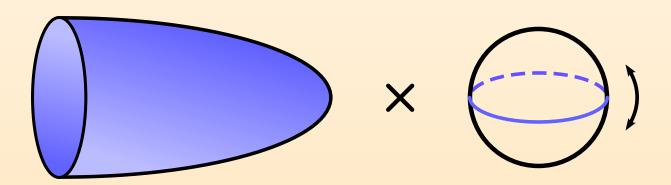
Karch, Katz 2002

	0	1	2	3	4	5	6	7	8	9
N D3	X	X	X	X						
1,2 D7	X	X	X	X	X	X	X	X		

Quarks: Low-energy limit of open strings between D3- and D7-branes

Meson masses from fluctuations of the D7-brane as given by DBI action:

Mateos, Myers, Kruczenski, Winters 2003

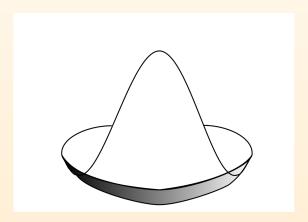


## Light mesons

Babington, J.E., Evans, Guralnik, Kirsch hep-th/0306018

Probe brane fluctuating in confining background:

Spontaneous breaking of  $U(1)_A$  symmetry



New ground state given by quark condensate  $\langle \bar{\psi}\psi \rangle$ 

Spontaneous symmetry breaking  $\rightarrow$  Goldstone bosons

Mass of  $\rho$  meson as function of  $\pi$  meson mass<sup>2</sup> (for  $N \to \infty$ )

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### Gauge/gravity duality:

 $\pi$  meson mass from fluctuations of D7-brane embedding coordinate

Bare quark mass determined by embedding boundary condition

 $\rho$  meson mass from D7-brane gauge field fluctuations

J.E., Evans, Kirsch, Threlfall 0711.4467

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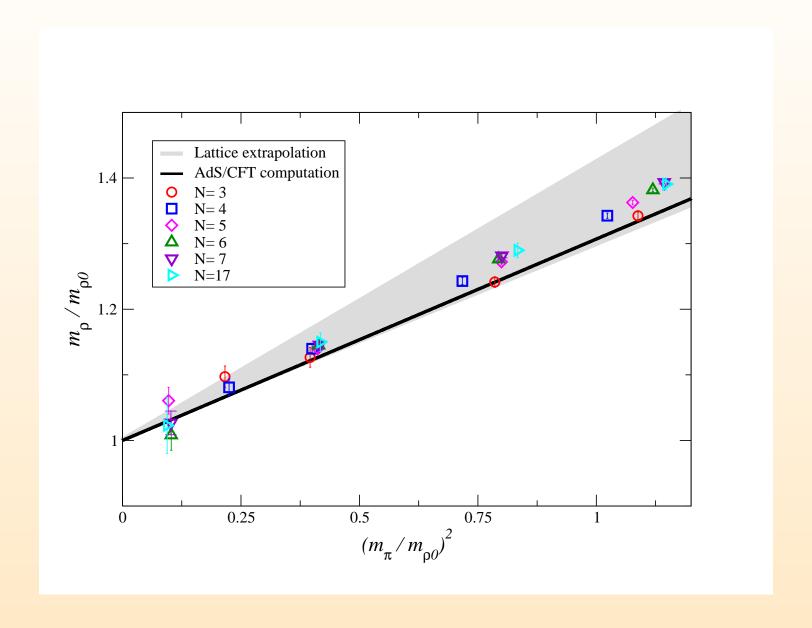
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Lattice: Bali, Bursa, Castagnini, Collins, Del Debbio, Lucini, Panero 1304.4437



D7 probe brane DBI action expanded to quadratic order:

$$S = \tau_7 \text{Vol}(S^3) \text{Tr} \int d^4x d\rho \, \rho^3 \left[ \frac{1}{\rho^2 + |X|^2} |DX|^2 + \frac{\Delta m^2 R^2}{\rho^2} |X|^2 + (2\pi\alpha' F)^2 \right]$$

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### Phenomenological model:

Evans, Tuominen 1307.4896

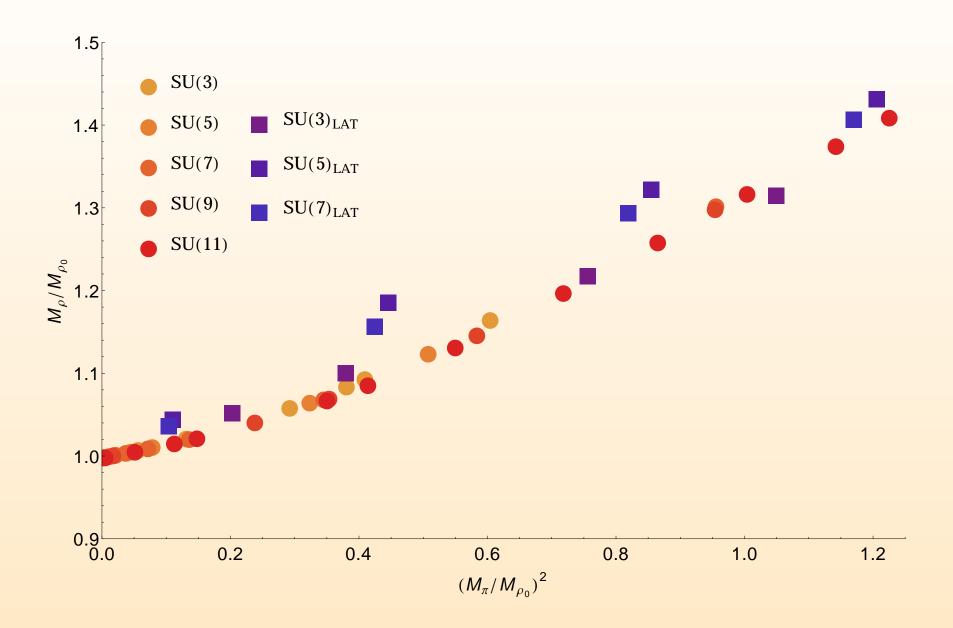
Metric

$$ds^{2} = \frac{R^{2}d\rho^{2}}{\rho^{2} + |X|^{2}} + \frac{(\rho^{2} + |X|^{2})}{R^{2}}dx^{2}$$

Fluctuations  $X = L(\rho)e^{2i\pi^a T^a}$ 

Make contact with QCD by chosing

$$\Delta m^2 R^2 = -2\gamma = -\frac{3(N^2 - 1)}{2N\pi} \alpha$$



### Bottom-up AdS/QCD model:

Chiral symmetry breaking from tachyon condensation

latrakis, Kiritsis, Paredes 1003.2377, 1010.1364

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latrakis, Kiritsis, Paredes 1003.2377, 1010.1364

### SU(N) Yang-Mills theory

Panero: Lattice studies of quark-gluon plasma thermodynamics 0907.3719

Pressure, stress tensor trace, energy and entropy density

Comparison with AdS/QCD model of Gürsoy, Kiritsis, Mazzanti, Nitti 0804.0899

### 4. Axial anomalies

J.E., Haack, Kaminski, Yarom 0809.2488;

Banerjee, Bhattacharya, Bhattacharyya, Dutta, Loganayagam, Surowka 0809.2596

Action of  $\mathcal{N}=2$ , d=5 Supergravity:

From compactification of d=11 supergravity on a Calabi-Yau manifold

$$S = -\frac{1}{16\pi G_5} \int \left[ \sqrt{-g} \left( R + 12 - \frac{1}{4} F^2 \right) - \frac{1}{2\sqrt{3}} A \wedge F \wedge F \right] d^5 x$$

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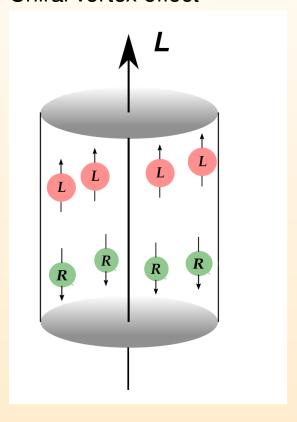
Contribution to relativistic hydrodynamics, proportional to angular momentum:

$$J_{\mu}=
ho u_{\mu}+\xi\omega_{\mu}, \quad \omega_{\mu}=rac{1}{2}\epsilon_{\mu
u\sigma
ho}u^{
u}\partial^{\sigma}u^{
ho}$$
, in fluid rest frame  $\vec{J}=rac{1}{2}\xi
abla imes \vec{v}$ 

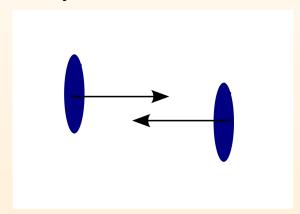
### Chiral vortex effect

Chiral separation: In a volume of rotating quark matter, quarks of opposite helicity move in opposite directions. (Son, Surowka 2009)

Chiral vortex effect



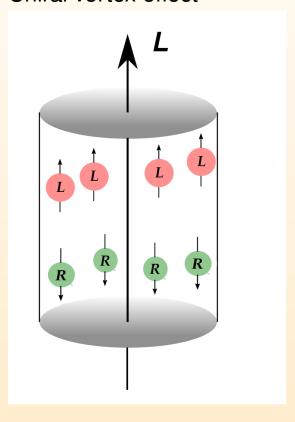
Non-central heavy ion collision



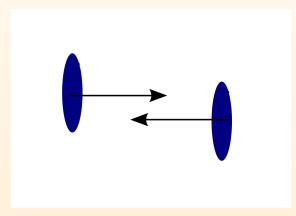
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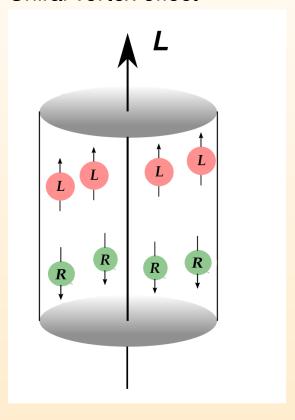


Chiral vortex effect ⇔ Chiral magnetic effect Kharzeev, Son 1010.0038; Kalaydzhyan, Kirsch 1102.4334

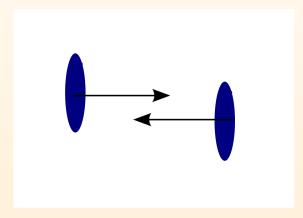
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Non-central heavy ion collision



Chiral vortex effect ⇔ Chiral magnetic effect Kharzeev, Son 1010.0038; Kalaydzhyan, Kirsch 1102.4334

Anomaly induces topological charge  $Q_5 \Rightarrow$  Axial chemical potential  $\mu_5 \leftrightarrow \Delta Q_5$  associated to the difference in number of left- and right-handed fermions



# Chiral vortex effect for gravitational axial anomaly

### Similar analysis for gravitational axial anomaly

$$\partial^{\mu} J_{\mu}^{5} = a(T) \, \varepsilon_{\mu\nu\rho\sigma} R^{\mu\nu}{}_{\alpha\beta} R^{\rho\sigma\alpha\beta}$$

Both holographic and field-theoretical analysis reveal  $a(T) \propto T^2$ 

Landsteiner, Megias, Melgar, Pena-Beñitez 1107.0368 Landsteiner, Megias, Pena-Beñitez 1103.5006 (QFT) Chapman, Neiman, Oz 1202.2469 Jensen, Loganayagam, Yarom 1207.5824

# Chiral vortex effect for gravitational axial anomaly

### Linear response

$$\vec{J}^5 = \sigma \vec{\omega}$$

$$\sigma = \lim_{p_j \to 0} \sum_{i,k} \frac{i}{p_j} \epsilon_{ijk} \langle J_5^i(\vec{p}) T^{k0}(0) \rangle \sim \frac{T^2}{24}$$

## Chiral vortex effect for gravitational axial anomaly

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Conversely,

$$J_E^i = T^{0i} = \sigma B_5^i$$

 $B_5$  axial magnetic field

couples with opposite signs to left-and right-handed fermions

Axial magnetic effect Braguta, Chernodub, Landsteiner, Polikarpov, Ulybyshev 1303.6266

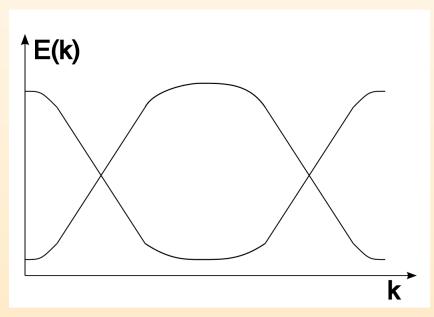
# Proposal for experimental observation in Weyl semimetals

Chernodub, Cortijo, Grushin, Landsteiner, Vozmediano 1311.0878

Semimetal: Valence and conduction bands meet at isolated points

Dirac points: Linear dispersion relation  $\omega = v|\vec{k}|$ , as for relativistic Dirac fermion

Weyl fermion: Two-component spinor with definite chirality (left- or right-handed)



Band structure of Weyl semimetal

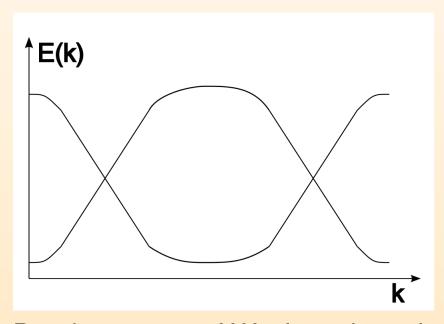
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Band structure of Weyl semimetal

Experimental observation of Dirac semimetals: '3D graphene'

Cd<sub>3</sub>AS<sub>2</sub>: 1309.7892 (Science),

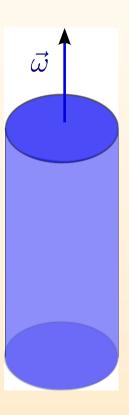
1309.7978

Na<sub>3</sub>Bi: 1310.0391 (Science)

## Proposal for experimental observation in Weyl semimetals

Chernodub, Cortijo, Grushin, Landsteiner, Vozmediano 1311.0878

- Weyl points separated by wave vector
- Wave vector corresponds to axial vector potential
- This induces an axial magnetic field at edges of a Weyl semimetal slab
- Via Kubo relation this generates angular momentum  $L_k = \int_V \epsilon_{ijk} \, x^i T^{0j}$
- By angular momentum conservation, this leads to a rotation of the slab
- This depends on  $T^2$



# **Summary**

- 1. Holographic Kondo model: RG flow
- 2. New inhomogeneous ground states
- 3. Mesons: Comparison to lattice gauge theory
- 4. Axial anomalies: Quark-gluon plasma ⇔ Condensed matter physics

### At this conference:

- Quantum phases of matter
- Time dependence (Turbulence, non-equilibrium, quantum quenches)
- Holographic entanglement entropy
- Lattices and transport

### **Conclusion and Outlook**

Gauge/gravity duality:
 Established approach for describing strongly coupled systems

### Conclusion and Outlook

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   Established approach for describing strongly coupled systems
- Unexpected relations between different branches of physics ⇔ Universality
- Comparison of results with lattice gauge theory, effective field theory, condensed matter physics
- Gauge/gravity duality has added a new dimension to string theory

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#### In the future:

- Mutual influence: Fundamental ⇔ applied aspects of gauge/gravity duality
- First step: 1/N,  $1/\sqrt{\lambda}$  corrections