#### Holographic Complexity Equals Which Action?

Leonel Queimada

King's College London

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 $9^{th}$  July 2019





#### Based on

Goto, Marrochio, Myers, LQ, Yoshida JHEP 1902 (2019) 160, arXiv:1901.00014

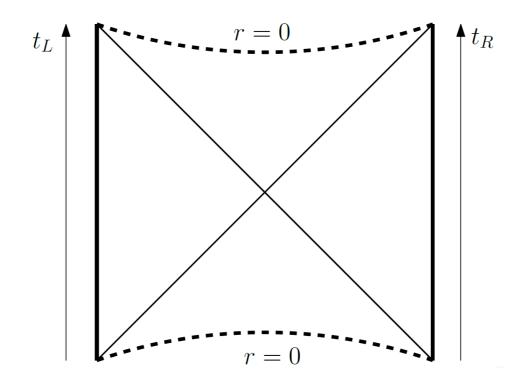
## See also

Brown, Gharibyan, Lin, Susskind, Thorlacius, Zhao Phys.Rev. D99 (2019) no.4, 046016, arXiv:1810.08741

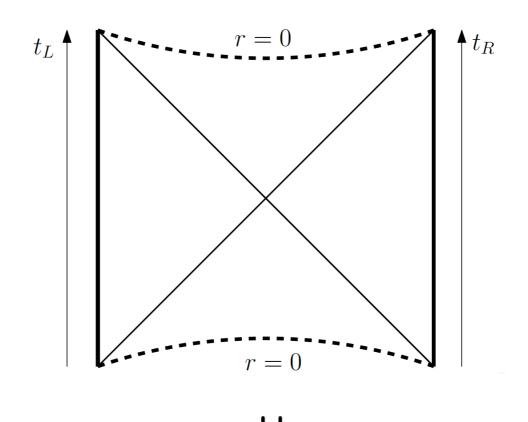
Alishahiha

Eur.Phys.J. C79 (2019) no.4, 365, arXiv:1811.09028

#### Eternal Black Hole = Thermofield Double State



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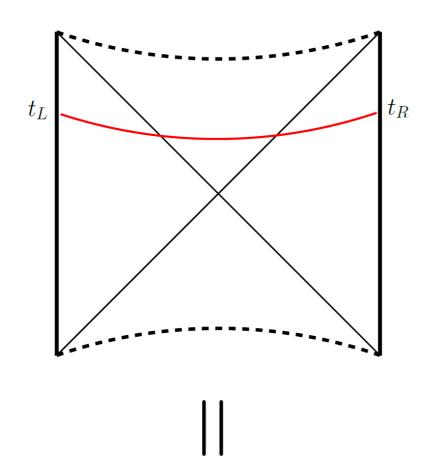


$$|TFD(t_L, t_R)\rangle = Z^{-1/2} \sum_{\alpha} e^{-E_{\alpha}/(2T)} e^{-iE_{\alpha}(t_L + t_R)} |E_{\alpha}\rangle_L |E_{\alpha}\rangle_R$$

(Maldacena, 2001)

(Hartman, Maldacena, 2013)

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(Susskind, Stanford, 2014)

### What is the Complexity of a State?

 $\mathcal{C}(|\Psi\rangle) = \text{minimum size of quantum circuit to prepare the state}$ 

Computer Science, Simple Quantum Systems

(Aaronson, 2016)

(Watrous, 2008)

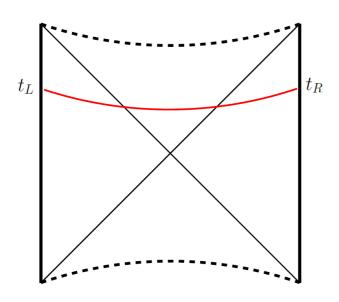
Quantum Field Theories

(Jefferson, Myers, 2017)

(Chapman et al, 2017)

(Caputa *et al*, 2017)

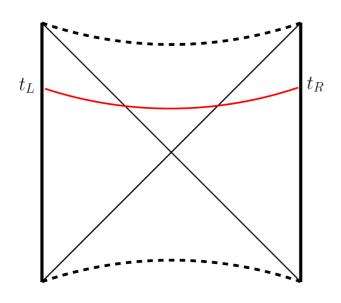
## Holographic Complexity Conjectures



$$C_{\mathcal{V}} = \max \left[ \frac{\mathcal{V}}{Gl} \right]$$

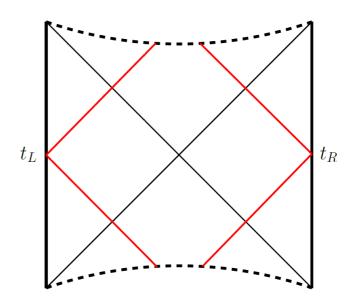
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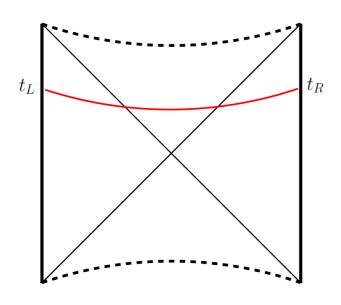
(Susskind, Stanford, 2014)



$$C_{\mathcal{A}} = \frac{I_{\text{WDW}}}{\pi \hbar}$$

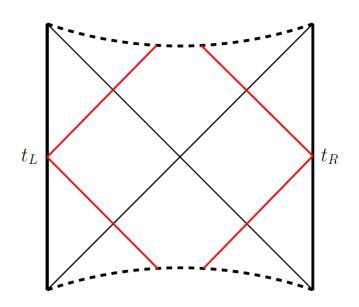
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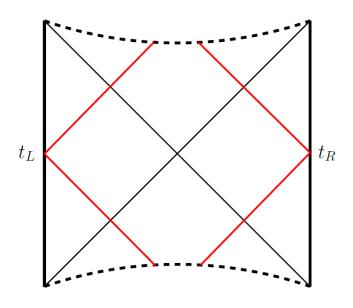


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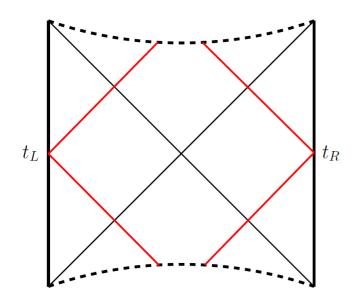
(Brown *et al*, 2016)

$$\lim_{t \to \infty} \frac{d\mathcal{C}}{dt} \sim TS$$

Let us first consider a Schwarzschild-AdS black hole



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$$I_{\rm WDW} = I_{\rm EH} + I_{\rm surf} + I_{\rm ct}$$

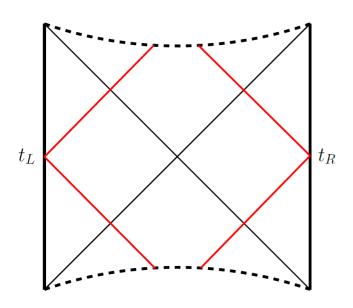
$$I_{\rm EH} = \frac{1}{16\pi G_N} \int_{\mathcal{M}} d^4x \sqrt{-g} \left(\mathcal{R} - 2\Lambda\right)$$

$$I_{\text{surf}} = \frac{1}{8\pi G_N} \int_{\mathcal{B}} d^3x \sqrt{|h|} K + \frac{1}{8\pi G_N} \int_{\Sigma} d^2x \sqrt{\sigma} \eta$$
$$+ \frac{1}{8\pi G_N} \int_{\mathcal{B}'} d\lambda \, d^2\theta \sqrt{\gamma} \kappa + \frac{1}{8\pi G_N} \int_{\Sigma'} d^2x \sqrt{\sigma} a$$

$$I_{\rm ct} = \frac{1}{8\pi G_N} \int_{\mathcal{B}'} d\lambda \, d^2\theta \sqrt{\gamma} \,\Theta \, \log \left(l_{ct}\Theta\right)$$

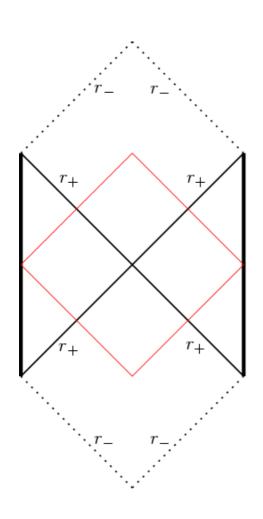
(Lehner et al, 2016)

Let us first consider a Schwarzschild-AdS black hole.

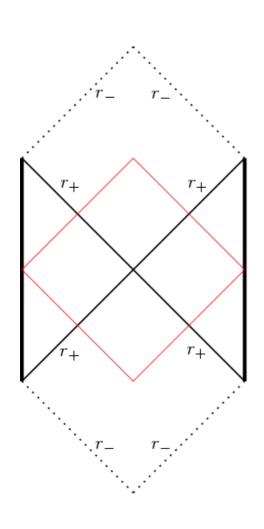


$$\lim_{t \to \infty} \frac{d\mathcal{C}_{\mathcal{A}}}{dt} = \frac{2M}{\pi\hbar}$$

Now, we consider a dyonic Reissner-Nordstrom-AdS black hole



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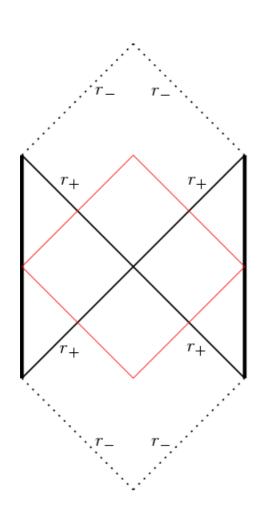


$$I_{\rm WDW} = I_{\rm Grav} + I_{\rm Max}$$

$$I_{\text{Max}} = -\frac{1}{4g^2} \int_{\mathcal{M}} d^4 x \sqrt{-g} \, F_{\mu\nu} F^{\mu\nu}$$

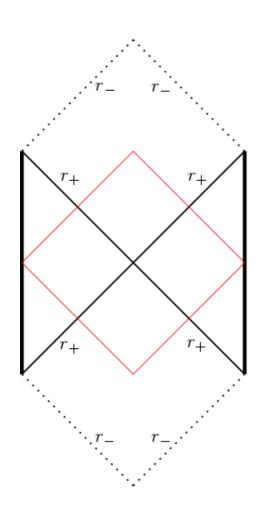
$$F = \frac{g}{\sqrt{4\pi G_N}} \left( \frac{q_e}{r^2} dr \wedge dt + q_m \sin\theta d\phi \wedge d\theta \right)$$

Now, we consider a dyonic Reissner-Nordstrom-AdS black hole.



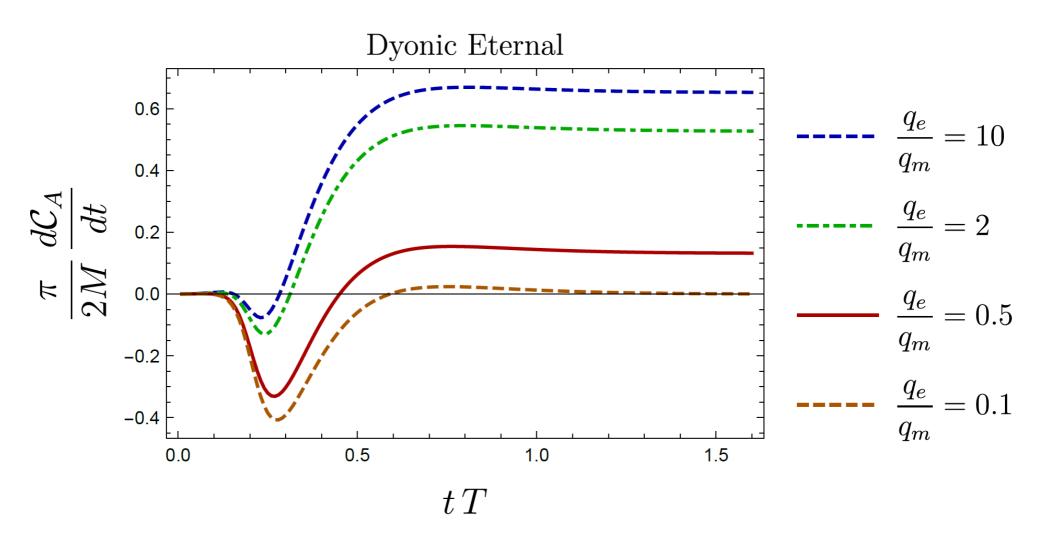
$$\lim_{t \to \infty} \frac{d\mathcal{C}_{\mathcal{A}}}{dt} = \frac{1}{\pi \hbar} \left[ \frac{q_e^2}{G_N r} \right]_{r_+}^{r_-}$$

Now, we consider a dyonic Reissner-Nordstrom-AdS black hole.



$$\lim_{t \to \infty} \frac{d\mathcal{C}_{\mathcal{A}}}{dt} = \frac{1}{\pi \hbar} \left[ \frac{q_e^2}{G_N r} \right]_{r_+}^{r_-}$$

The magnetic charge seems to play no role.

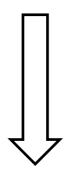


#### Lessons from BH thermodynamics

For simplicity, let us assume a purely electric black hole  $(q_m = 0)$ .

$$I_{\text{Eucl}} = I_{\text{Grav}} + I_{\text{Max}}$$

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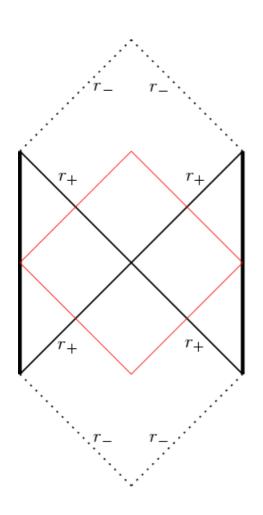
$$\int \int I_{\mu Q} = \frac{1}{g^2} \int_{\partial \mathcal{M}} d\Sigma_{\mu} F^{\mu\nu} A_{\nu}$$

Fixed  $\mu$  ensemble

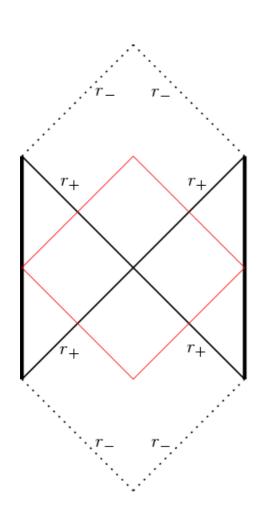
Fixed Q ensemble

(Hawking, Ross, 1995)

We can consider the addition of a boundary term for the Maxwell field



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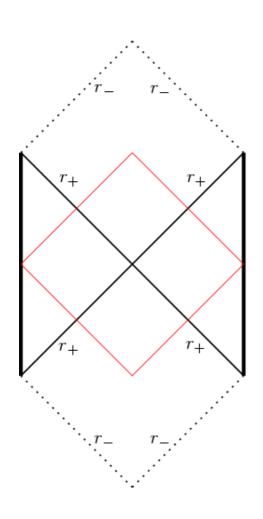


$$I_{\text{WDW}} = I_{\text{Grav}} + I_{\text{Max}} + I_{\mu Q}$$

$$I_{\mu Q} = \frac{1}{g^2} \int_{\partial \mathcal{M}} d\Sigma_{\mu} F^{\mu\nu} A_{\nu}$$

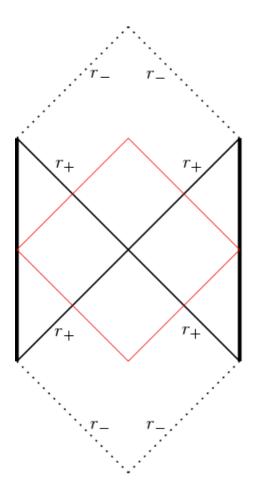
$$F = \frac{g}{\sqrt{4\pi G_N}} \left( \frac{q_e}{r^2} dr \wedge dt + q_m \sin\theta \, d\phi \wedge d\theta \right)$$

We can consider the addition of a boundary term for the Maxwell field



$$\lim_{t \to \infty} \frac{d\mathcal{C}_{\mathcal{A}}}{dt} = \frac{1}{\pi \hbar} \left[ \frac{q_m^2}{G_N r} \right]_{r_+}^{r_-}$$

We can consider the addition of a boundary term for the Maxwell field



$$\lim_{t \to \infty} \frac{d\mathcal{C}_{\mathcal{A}}}{dt} = \frac{1}{\pi \hbar} \left[ \frac{q_m^2}{G_N r} \right]_{r_+}^{r_-}$$

The boundary term changes the type of charge that contributes to the complexity growth rate.

$$I_{\rm BH} = I_{\rm Grav} + I_{\rm Max}$$

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$$ds^{2} = g_{ab}dx^{a}dx^{b} + r^{2}d\Omega^{2}$$

$$F = \frac{g}{\sqrt{4\pi G_{N}}}q_{m}\sin\theta \,d\phi \wedge d\theta$$

$$\Phi = 4\pi r^{2}$$

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$$I_{\rm JT} = \frac{\Phi_0}{16\pi G_N} \left( \int_M d^2x \sqrt{-g}R + 2 \int_{\partial M} \sqrt{|\gamma|}K \right) + \frac{1}{16\pi G_N} \left( \int_M d^2x \sqrt{-g}\Phi_{\epsilon} \left( R + \frac{2}{L_2^2} \right) + 2 \int_{\partial M} \sqrt{|\gamma|}\Phi_{\epsilon}K \right)$$

(Nayak *et al*, 2018)

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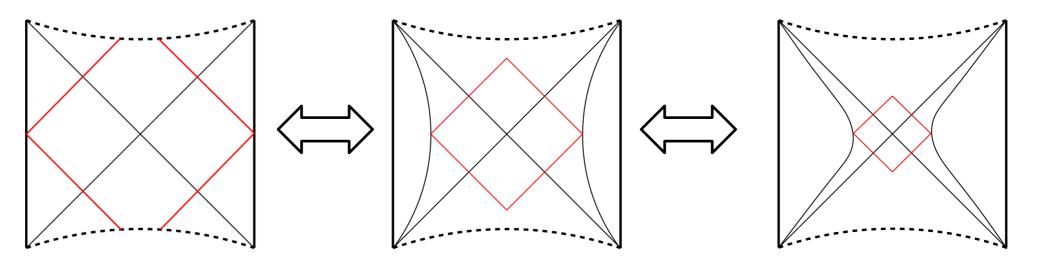
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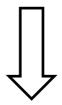
(Nayak *et al*, 2018)



JT model action does not grow at late times.

#### Main Results and Outlook

- o Late time complexity growth is independent of where we anchor the WDW patch. Possible relation with  $T\bar{T}$  deformations? (Akhavan *et al*, 2018)
- O Boundary Term Dependence of Holographic Complexity



Distinction between CA and CV conjectures

- What is the meaning of the boundary terms associated to the matter fields for the definition of complexity in the dual CFT?
- More generally, given the ambiguities associated to an action, holographic complexity equals which action?